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11. 1.

PERFORMANCE ANALYSIS OF A ROTARY REGENERATOR

A THESIS

Presented to

The Faculty of the Graduate Division

by

Humberto Eduardo Barrientos-Mendoza

In Partial Fulfillment

of the Requirements for the Degree

Master of Science

in Mechanical Engineering

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March, 1971

PERFORMANCE ANALYSIS OF A ROTARY REGENERATOR

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NOMENCLATURE

vii

surface area, ft²

f

h

k

L

Ð

N

NTŰ

NTUo

P

p

Q

r

T

ν

v

W

RHE

specific heat, Btu/1bm °R coefficient of heat transfer by convection, Btu/hr ft² °F specific enthalpy of air, Btu/1bm thermal conductivity of the matrix, Btu/hr ft °F length of the matrix, ft mass rate flow, 1bm/hr

number of elements of matrix number of heat transfer units modified number of heat transfer units =

 $\frac{1}{C_{\min}} \left[\frac{1}{(1/hA)_{c} + (1/hA)_{h}} \right]$

pressure, 1bf/ft²

free flow area ratio of the matrix =

free flow cross-sectional area total cross-sectional area

heat transfer rate, Btu/hr

rotary heat exchanger

radial coordinate in the cylindrical frame of reference temperature, °F

specific volume, ft³/1bm

velocity, ft/hr

specific humidity, 1bm of water/1bm dry air

NOMENCLATURE (Concluded)

axial coordinate in the cylindrical frame of reference

hot or cold

GREEK LETTERS

z

x

Ø

α,

Ø

ρ

a

avg

С

đ

е

h

i

ŕ

rw

s

t

17

æ

angular velocity, rad/hr

surface area density, ratio of total surface area of the matrix to the volume of the matrix, 1/ft

circumferential coordinate in the cylindrical frame of reference

density, 1bm/ft³

SUBSCRIPTS

air

average

cold

conduction

external

hot

internal

rotor or matrix

effective average over rotor and condensed heater

at the surface of the matrix

total

convection

water

free stream condition

SUMMARY

This thesis concerns the application of a numerical technique to study the performance characteristics of a periodic-flow heat exchanger, with at least one fluid having sufficiently high relative humidity to cause condensation.

Suitable analytical and numerical models are presented to describe the heat exchange between the fluids and the exchanger. Distribution of temperatures within the matrix of the heat exchanger and thermodynamic properties of the fluids (air) flowing through the exchanger are discussed.

CHAPTER I

INTRODUCTION

Rotary Regenerator Heat Exchanger

A commonly accepted classification of heat exchangers is made on the basis of whether they employ heat storage or not. Commonly termed a regenerator, the former operates by means of a storing material which runs through two heat exchanging media, absorbing heat from the hotter medium and transferring heat to the cooler one, through which it flows. On the other hand, a recuperator is a heat exchanger which does not employ heat storage. It operates by direct heat transmission through a separating wall between two fluids exchanging heat or by two direct heat transfer units, coupled with a transfer fluid circulating in cycles between the hot exchanger unit where the thermal energy is received and the cold exchanger unit where the thermal energy is delivered. Such an exchanger necessitates a simultaneous flow of both the fluids engaged in the mutual heat transfer process.

During the past three or four decades several designs incorporating a rotary matrix for the flow of material through which, in turn, flow two streams of fluids at different temperatures, have been con-

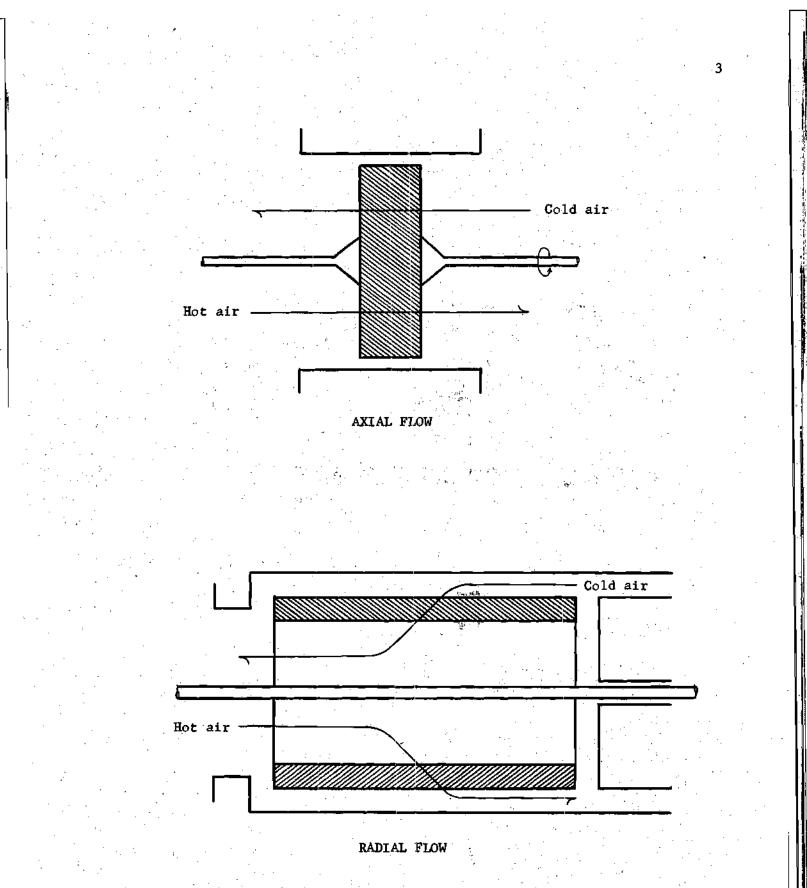
sidered. In the past, a primary motivation for heat exchangers of such a design, suitably termed as rotary regenerators, has been their use in

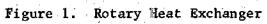
gas turbine power systems.

A further classification of rotary regenerators or rotary heat exchangers (RHE) can be made on the basis of whether the flow is axial or radial. Figure 1 depicts schematic representations of the two types of RHE.

A common feature of both of these RHE involves the rotation of a porous matrix from a cold fluid stream into a hot fluid stream in a regular periodic fashion. During the part of the cycle that the metal matrix interacts with the cold stream, the matrix loses part of its internal energy due to the heat transferred from the metal to the cold fluid stream. On a subsequent part of the cycle, when due to its rotation the matrix interacts with the hot fluid stream, the metal regains the internal energy via the heat transferred from the fluid to the matrix. For a given speed of rotation of the matrix, together with the proper mass fluxes of cold and hot streams, it is possible to achieve a steady state periodicity of temperature distribution within the matrix. In other words, at the end of a given cycle, the whole matrix attains the same state as it possessed at the end of the previous cycle. The net effect of such a cycle of operation is the exchange of heat between the hot and the cold fluid streams. In principle, this facilitates the cooling of a warm stream or the heating of a cold one, whichever serves the purpose in a particular application.

A perspective view of a typical matrix arrangement in the axial flow type of RHE is shown in Figure 2. Two advantageous features of such a design concern its compactness and the tendency of self cleaning. The heat transfer area per unit volume or the surface area density in such an





Ŧ.

Hot air in

Figure 2. Matrix of a Rotary Heat Exchanger

10

Cold air out

exchanger is higher than that in a conventional exchanger, thus requiring less volume for a given heat transfer area; this is related to the structure of the porous matrix. Since the axial direction of flow through a particular sector of the matrix is reversed during each cycle, such a geometry provides the self-cleaning feature.

A Brief Survey of Pertinent Literature

Well through the 1930's, two procedures were employed to predict the performance of regenerators (1). One of these procedures was to use empirical data to set up approximate equations by means of which the performance of the regenerator could be predicted. However, because of the nature of the empirical analysis, these equations were only fair approximations for those cases similar to the ones for which the empirical data were obtained. Naturally, they could not be used for other specific cases. The other procedure was based upon analytically or numerically derived solutions of the problem. This last method proved to be very successful and it is being adopted even at present.

Hausen (2) analyzed the problem of RHE by the so-called heat-pole method. He considered the thermal conductivity of the metal of a cylindrical shaped matrix negligible in the direction of air flow (axis of cylinder) and to be infinite in the radial direction. He also considered the heat capacity rate (product of mass flow rate and specific heat) of the hot air flowing through the RHE to be the same as that for the cold matrix itself. He further assumed a unity value for the ratio of the heat conductance of the cold side to the heat conductance of the hot side. By considering the matrix to be divided into an equal number of strips,

-

he obtained the dimensionless temperature distribution, as well as the total heat transfer rate and the regenerator efficiency. He found that the solution converged very quickly with an increase in the number of strips, namely 0.1 percent deviation between calculated efficiencies for nine and for ten strips.

In the 1940's the aircraft industry showed a great interest in the rotary regenerator as a direct application to improving the performance of gas turbines. Substantial attention was focused on the effect of leakage on regenerator effectiveness, and a consequent development of satisfactory seals to minimize leakage, and on the role played by such parameters as the matrix dimensions, as well as the velocities and heat transfer coefficients associated with the problem.

Harper and Rohsenow (3) showed that, in gas turbine applications of the RHE, most of the leakage occurred not by carry-over of the matrix but through the seals of the matrix from the high pressure stream to the low pressure one. In their paper the change of efficiency of the regenerator was evaluated, by considering a relative change of temperature of the discharging stream due to leakage as compared with that for the case of no leakage. This change of efficiency was shown to be rather small, about 1.3 percent corresponding to a 10 percent leakage. They also showed that the effectiveness of the regenerator appeared to increase to a maximum with an increase in thickness of the matrix, beyond which the thermodynamic efficiency of the gas turbine cycle was observed to drop due to an increase in loss of pressure.

The influence of gas stream velocities and mass flow rates on the

heat transfer characteristics of the matrix was determined by Tong and London (4) for various matrix structures. Their work was primarily of an empirical nature, with the results expressed as plots of four dimensionless parameters, namely, the drag coefficient, the friction factor, the Reynolds number, and the Prandtl number.

The search for analytical means to predict the performance of the periodic flow regenerator was continued by Coppage and London (5) in a paper where they first made a survey of the different available solutions for the design of regenerators, and then developed a closed form approximate solution of their own which incorporated the work of previous investigators.

Dusinberre (5), in his discussion of the paper by Coppage and London, suggested that, by dividing the matrix of a regenerator into a number of small elements, each element being treated as a cross flow heat exchanger, the efficiency of the regenerator could be calculated for a given set of physically meaningful parameters. Indeed, this was a very useful and powerful suggestion. Lambertson (6) executed Dusinberre's suggestion and developed a finite difference numerical technique, thus contributing significantly toward the solution of the problem, which had eluded many previous investigators. His solutions were useful for a variety of gas turbine applications.

Lambertson's work forms a commonly accepted basis for design data for the periodic flow type of heat exchanger and has been incorporated by Kays and London (7) in their textbook on compact heat exchangers. It is worth noting that the influence of axial conduction had not been included in Lambertson's work. It was generally accepted that axial conduction of heat might be a significant consideration for high performance regenerators.

Bahnke and Howard (8) carried out Dusinberre's suggestion in a fashion similar to that by Lambertson, but they included the effect of longitudinal heat conduction in the heat exchanger. It was shown in their paper that, for a regenerator effectiveness lower than 90 percent, axial conduction did not play a significant role in the overall heat transfer mechanism. However, they showed that, for effectiveness higher than 90 percent, the influence of axial heat conduction was progressively important in the computation of effectiveness and heat transfer units.

Recent years have seen a considerable increase of interest in utilizing the rotary regenerators beyond the classical gas turbine applications. Drying and air conditioning systems are but two of the more prominent cases where the use of rotary regenerators may show considerable promise. In such applications, the hot and cold streams of air, owing to their thermodynamic states, involve condensation of the water vapor flowing with the air. In an attempt to establish the applicability of a rotary regenerator to a clothes drying process, Mercure (9) performed simple thermodynamic computations to evaluate its performance. In such an application, room air sucked through a sector of the regenerator matrix and passed through an electric heater is supplied to the drum of the clothes dryer. After absorbing moisture from a wet load of clothes, the hot, humid air is passed through another sector of the matrix, and is subsequently exhausted into the ambient. The rotary matrix absorbs heat from the hot stream, stores the energy, and returns it to the incoming cold stream in a periodic fashion. In the process, a certain amount of condensate is undesirably carried over from the hot sector to the cold one. Due to its simplistic nature, Mercure's computational procedure did not allow a prediction of regenerator effectiveness based on simultaneous condensation and convective heat transfer. No other references to the analyses of a rotary heat exchanger in the presence of condensing water vapor on the matrix are available in the literature at present.

Statement of the Proposed Problem

The object of this study is to investigate the application of a finite difference numerical technique to the performance analysis of an RHE with humid air flowing through it. The approach adopted for this investigation consists of developing a technique to predict the variable humid air properties leaving the regenerator for a given state of entering air, with the results coupled to an effectiveness analysis capable of predicting the performance of the rotary regenerator.

CHAPTER II

DEVELOPMENT OF THE ANALYTICAL MODEL

Selection of the Coordinate System and the Differential

Control Volume

The physical geometry of the domain of interest in this analysis, namely a cylindrical matrix of a rotary heat exchanger, suggests the use of a cylindrical coordinate frame. The origin is taken at the center of the top of the matrix. The positive z direction is measured downward along the axis of symmetry. The variable r is measured outward along the radius of the cylinder, and \emptyset is measured positive in the counterclockwise direction. With reference to such a coordinate frame, the domain of the matrix can be defined as:

 $r_{i} \leq r \leq r_{e},$ $0 \leq \emptyset \leq 2\pi,$ $0 \leq z \leq L.$

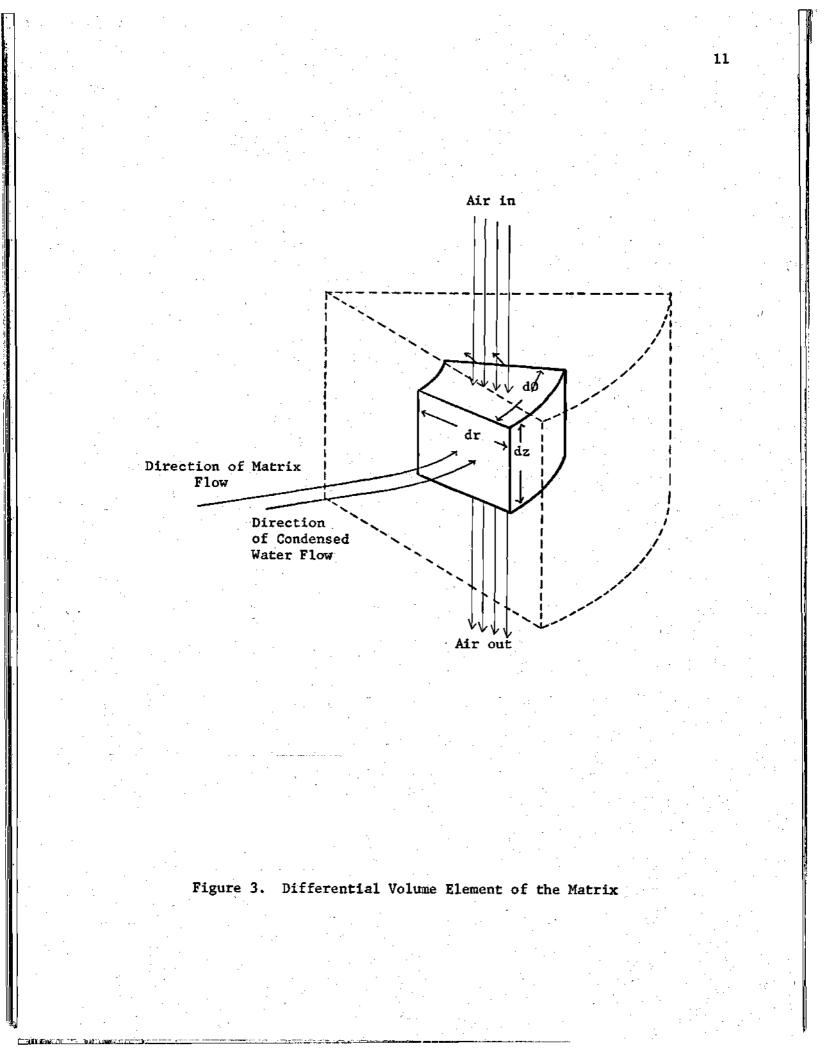
A differential control volume fixed in space and referred to the cylindrical coordinate frame is shown in Figure 3.

As the matrix rotates, its elements may be considered to enter and leave the differential control volume in a periodic fashion. The control volume is considered to be interacting with its surroundings in

the following fashion:

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(2-1)



Matrix and water enter and leave the control volume along the Ø direction at constant velocity, thus resulting in transport of heat and mass across the control surface. Humid, hot air enters through the top and leaves through the bottom of the control volume along the positive z direction at constant flow rate, with mass and energy transport taking place, as explained above. Heat transfer by conduction enters the three positive faces and leaves through the three opposite faces of the control surface.

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(2-2)

Governing Equations and Boundary Conditions

A steady state heat balance over the control volume may be expressed, in words, as

> (Net conduction heat transfer into the C.V.) + (Net heat transfer in, due to the matrix and water mass entering and leaving the C.V.) + (Net heat transfer due to the air entering the top and leaving through the bottom of the C.V.) = 0.

This same equation, written on a per unit volume basis in terms of symbols, gives

 $\left[k_{r}\left(\frac{\partial^{2}T}{\partial r^{2}}+\frac{1}{r}\frac{\partial T}{\partial r}\right)+\frac{k_{p}}{r^{2}}\frac{\partial^{3}T}{\partial \phi^{2}}+k_{z}\frac{\partial^{2}T}{\partial z^{2}}\right](1-p) - (2-3)$

 $\frac{\partial}{\partial \phi} (p \rho_{rw} C_{prw} \dot{\phi}T) + p \rho_{a} V_{a} \frac{\partial h_{\infty}}{\partial z} = 0 .$

Following Stoecker (10), the differential heat transfer from the humid

air to the C.V. is

$$d\dot{Q}_a = \frac{f}{C_{pm}} (h_{\infty} - h_s) dA$$
, (2-4)

$$d\dot{Q}_{a} = p \rho_{a} V_{a} \frac{\partial h_{\infty}}{\partial z} rd\phi drdz . \qquad (2-5)$$

The differential surface area available for convective heat trans-

$$A = \alpha r \phi dr dz$$
 .

where α is the surface area density, or the ratio of total surface area of the matrix to the volume of the matrix.

When equations (2-4), (2-5), and (2-6) are combined, one obtains

$$\frac{f}{C_{pm}} (h_{\infty} - h_{s}) dA = p \rho_{a} V_{a} \frac{\partial h_{\infty}}{\partial z} \frac{dA}{\alpha} . \qquad (2-7)$$

The boundary condition for enthalpy may be specified as

$$\begin{array}{c} 0 \leq \phi \leq \phi_{h} \\ z = 0 \end{array} \right\} \quad h(r, \phi) = \text{constant} \quad .$$
 (2-8)

The interval of existence of \emptyset in equation (2-8) specifies the hot sector side of the matrix.

In a similar fashion, for the cold side of the matrix, the boundary

(2-6)

condition may be written as

 $h(r, \emptyset, L) = constant.$

In these equations,

$$T = T(r, \phi, z), h_{\infty} = h_{\infty}(T, P_{a}), \text{ and } \rho_{a} = \rho_{a}(T, P_{a}).$$
 (2-10)

It may be pointed out that the summation of partial pressures of air and water yields the constant value of the total pressure, P_+ .

Equations (2-3) together with the differential energy balances stated in equations (2-4) and (2-5) as well as the boundary conditions expressed by equations (2-8) and (2-9) represent an analytical model of the problem.

<u>Simplification of the Analytical Model</u>

A closed form analytical solution of the preceding system of equations proves to be a formidable task. It is recognized that enthalpy and density are not known <u>a priori</u> as a function of spatial location; moreover, if spatial location is not involved, these properties are nonlinear functions of temperature and pressure.

These equations are elliptic-type partial differential equations. Classical mathematical techniques have been often employed with special treatment of this type of equation. However, general solutions for elliptic-type partial differential equations are not available.

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(2-9)

It is fortunate that some reasonable idealizations that do not significantly affect the engineering applicability of the equations can be made to simplify the problem. Following such simplifications, it is feasible to develop a finite difference model of the equations, which can be solved by using numerical solution techniques.

Consider the following idealizations with subsequent justifications of the same:

1) The two fluid streams through the matrix represent a counterflow energy exchange system.

2) The thermodynamic properties of the entering fluids are uniform over the flow inlet cross sections and these properties remain constant with time.

 The convection heat transfer coefficients are constant along a flow line.

4) The thermal properties of the matrix are invariant with changes of temperature and time.

5) No mixing of the hot and cold fluid streams occurs, either as a result of direct leakage or due to carry over.

6) Regular periodicity exists for all properties within the ma-

trix elements, yielding quasi-steady conditions.

7) The matrix rotates at a constant angular velocity.

8) The water condensed at a particular station within the matrix remains there until reevaporated and the temperature of the condensed water is assumed to be equal to the temperature of the matrix in which the water rests. 9) The thermal conductivity of the porous matrix is negligibly small in the circumferential direction and extremely large in the radial direction. Thus, heat by conduction in the circumferential direction may be neglected, and the temperature along a radius may be assumed constant.

Since the first three assumptions are very commonly used and justified in the literature on the subject, it is not necessary to elaborate on those any further.

The fourth assumption implies that the density, specific heat, and other thermal properties of the metal matrix are independent of temperature. This is reasonable to accept for the range of temperatures in such applications as drying and air conditioning processes considered in this investigation.

The fifth assumption is well approximated when the width of the matrix is small, as is usually encountered in RHE to minimize pressure losses; or if the air velocity is high compared with the velocity of rotation of the matrix, the approximation represents the physical picture reasonably well. For example, with an air velocity of 700 ft/min through a one inch thick matrix turning at 15 rpm, the dead angle is around 0.64 degrees, which approximately equals 0.4 percent of one turn of rotation. Moreover, it must be recognized that, in a drying process, both flow streams have about the same total pressure, rendering a low driving potential for leakage flow. This, of course, is not the case in gas turbine applications, where the pressure ratio is quite high. The sixth assumption corresponds to the quasi-steady state of the

system. The transient state at the start of operation tends to reach a

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quasi-steady state as the number of cycles of operation increases. Ιń other words, after a certain number of cycles, the system, from an engineering point of view, becomes one in steady state, provided that there is no storage of water condensed from the hot fluid stream on the matrix. For the case with condensation, the following sequence of operation results in a reevaporation of the condensate: Water collected on the matrix is carried along by the matrix to a region where the temperature is higher. The specific humidity of air in contact with the condensate in this region is higher than the specific humidity of free stream air in the region. As a result, a mechanism similar to the evaporation of water from the wick around a wet bulb occurs causing the evaporation. If the condensed water does not reevaporate during one complete turn, there will be successive accumulations of condensed water during each turn, still allowing an eventual steady state system with water in the liquid phase being carried over beyond the matrix by the air stream. Such an amount of condensation need not be anticipated for the case under study.

The ninth idealization is based on the relatively small order of magnitude of circumferential temperature gradient compared with that in the axial direction. The thermal conductivity is assumed to be very large in the radial direction, since with all the other idealizations, the matrix elements along a particular radius have the same governing equations, the same boundary conditions, and the same time of exposure to both hot and cold streams. This results in a constant temperature along a radius. The matrix thus behaves as if the thermal conductivity along a

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radius were infinite.

A Note on the Qualitative Temperature Distribution

Within the Matrix

During the hot cycle a sector of the matrix raises its temperature at a faster rate initially and at a slower rate toward the end, the latter caused by a diminishing temperature differential between the matrix and the hot stream of fluid. In a similar fashion during the cold cycle the matrix lowers its temperature at a faster rate initially and at a slower rate toward the end.

The temperature distribution within the matrix at the beginning of the hot cycle is the same as that at the end of the cold cycle, since the spatial positions of the matrix for the two cases are identical. This is usually referred to as the reversal condition. Finally, the mean temperature of the matrix lies somewhere between the mean entering temperatures of the hot and cold streams.

Simplified Analytical Model

With the idealizations that have been made, the following simplifications are possible.

Idealization nine, namely that the thermal conductivity of the matrix in the circumferential direction is negligibly small, implies that the heat transfer by conduction is negligible and that

 $\frac{r^3}{r^3} \frac{\partial \phi^2}{\partial^2} = \theta .$

(2-11)

Similarly, if the temperature along a radius may be assumed to be constant, one obtains

$$\frac{\partial T}{\partial r} = \frac{\partial^2 T}{\partial r^2} = k_r \left(\frac{\partial^2 T}{\partial r^2} + \frac{1}{r} \frac{\partial T}{\partial r} \right) = 0 . \qquad (2-12)$$

From idealization number seven one obtains a constant angular velocity of the matrix, i.e., ϕ = constant.

When these assumptions are incorporated into equation (2-3), the resulting equation becomes

$$k_{z}(1 - p) \frac{\partial^{a} T}{\partial z^{2}} - \frac{\partial}{\partial \phi} (p \rho_{rw} C_{prw} \dot{\phi} T) + p \rho_{a} V_{a} \frac{\partial n_{\infty}}{\partial z} = 0.$$
 (2-13)

The boundary conditions for enthalpy expressed by equations (2-8)and (2-9) remain unchanged. Thus, the analytical formulation of the transport of mass and heat in a rotary regenerator with condensing vapor is represented by equations (2-7) and (2-13), together with the boundary conditions expressed by equations (2-8) and (2-9).

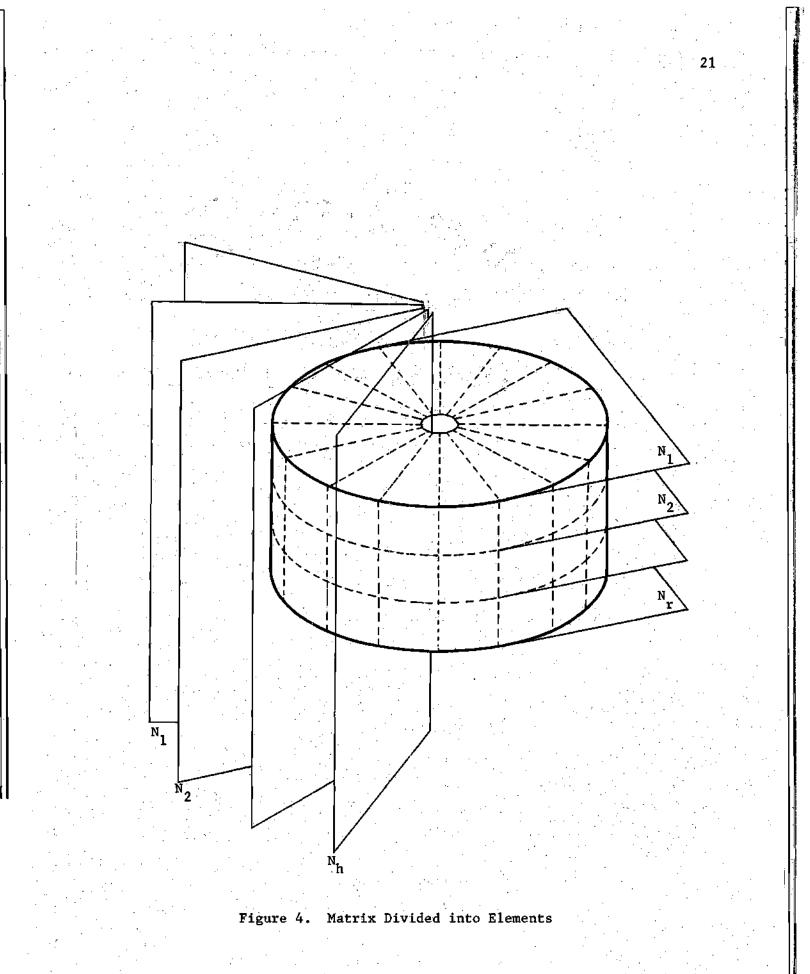
CHAPTER III

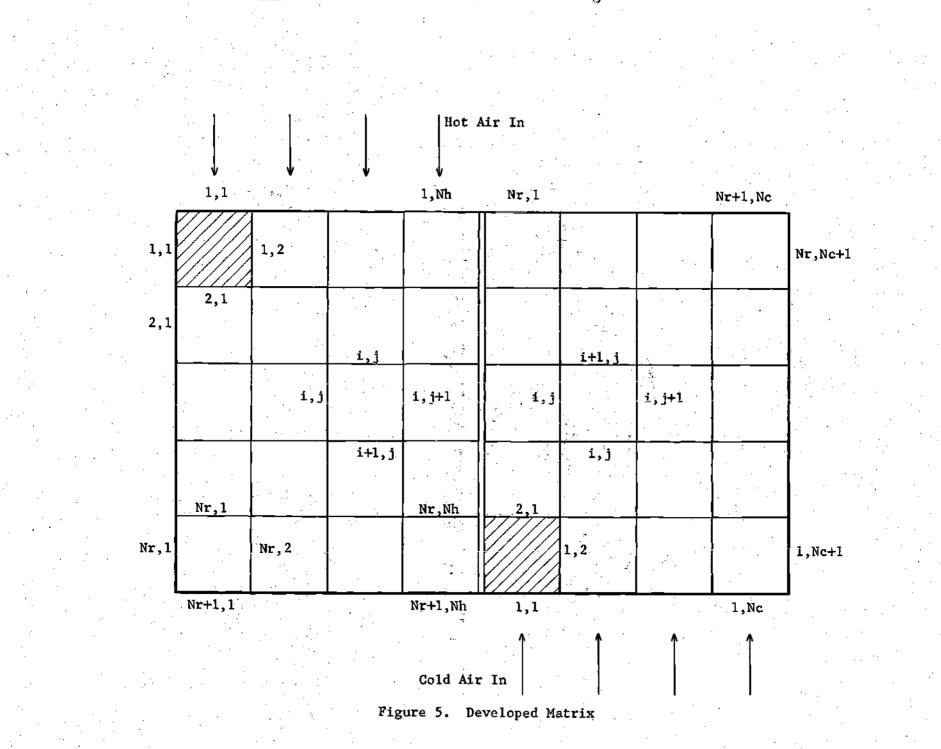
DEVELOPMENT OF THE NUMERICAL MODEL

The Grid System

For the purpose of developing a scheme of difference equations, the matrix is divided into N_r axial planes and $(N_h + N_c)$ radial planes around the circumference of the matrix. A total number of $N_r(N_h + N_c)$ elements, each element resembling the differential control volume discussed in Chapter II, are thus formed. It may be pointed out that, by distinguishing between the hot and cold sides of the matrix and by opening or developing the circular geometry onto a plane, the numerical analysis procedure can be carried out in the classical manner. The overall grid configuration bears a resemblance to a layered cake, as shown in Figure 4. The developed view of the grid system is shown in Figure 5.

The temperatures on the left and on the right sides of each element of the matrix are defined by a row-column system, the generic subscripts being i for rows and j for columns. The same generic system of description is used for the properties and states of the humid air at the top and bottom of each element. In other words, the generic subscript i implies a row and the generic subscript j implies a column. This scheme avoids any confusion and is simple because any reference to temperatures implies the left-right sides of the element, whereas all air states and properties refer to the top-bottom sides of the element.





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Finite Difference Approximation of the Analytical Model

It was pointed out in Chapter II that even the simplified analytical model leads to a system of differential equations for which a closed form solution is not available. A finite difference approximation of the model is therefore deemed necessary.

To proceed along these lines, equation (2-13) is first multiplied by the differential volume rdødrdz to yield

$$k_{z}(1 - p) \frac{\partial^{2} T}{\partial z^{2}} rd\phi drdz - \frac{\partial}{\partial \phi} (p \rho_{rw} C_{prw} \dot{\phi}T) rd\phi drdz - (3-1)$$

 $p \rho_a V_a \frac{\partial u_{\infty}}{\partial z} rd\phi drdz = 0 ,$

It was pointed out that (1 - p) rdødr is the area of the differential element available for conduction, which is denoted by dA_d . The factor (1 - p) is included to account for the porosity of the material of the matrix. pdrdz is the area through which water and matrix enter the differential volume. prdødr is the area through which air enters the differential volume.

The mass flow rate for air is given by

$$\dot{\mathbf{m}}_{a} = \rho_{a} \mathbf{V}_{a} \mathbf{A} = \rho_{a} \mathbf{V}_{a} \mathbf{prd} \mathbf{\phi} d\mathbf{r} , \qquad (3-2)$$

and the mass flow rate of the matrix element is given by

$$= \rho_V V_A = \rho_V p dr dz ,$$

A combination of equations (3-3) and (3-4) gives

where

 $V_r = r\phi$

$$\dot{m}_{r} = \rho_{r} r \phi p dr dz . \qquad (3-5)$$

The first term of equation (3-1), $k_z \frac{\partial^2 T}{\partial z^2} [(1-p) r d \phi d r] dz$, may be replaced by its equivalent $k_z \frac{\partial^2 T}{\partial z^2} dA_d dz$, or

$$k_{z} \frac{\partial^{2} T}{\partial z^{2}} [(1-p) r d \phi dr] \equiv k_{z} \frac{\partial^{2} T}{\partial z^{2}} dA_{d} dz . \qquad (3-6)$$

Contributions from water flow rate and matrix flow rate to the second term of equation (3-1) may be isolated to result in

$$\frac{\partial}{\partial \phi} \left[\rho_{\mathbf{rw}} \mathbf{C}_{\mathbf{prw}} \dot{\phi} \mathbf{T} \right] \mathbf{prd\phi} d\mathbf{r} dz \equiv \dot{\mathbf{m}}_{\mathbf{r}} \mathbf{C}_{\mathbf{pr}} \frac{\partial \mathbf{T}}{\partial \phi} d\phi + \mathbf{C}_{\mathbf{pw}} \frac{\partial (\dot{\mathbf{m}}_{\mathbf{w}} \mathbf{T})}{\partial \phi} d\phi \cdot (3-7)$$

It may be noted that the mass flow rate of water depends on an increase or decrease in the rate of condensation or evaporation. On the other hand, the mass flow rate of rotor is a constant quantity.

The third term of equation (3-1) may be expressed with the use of equation (3-2) as

$$p_{a}^{\partial h} = \frac{\partial h_{\infty}}{\partial z} rd\theta dr dz \equiv m_{a}^{\partial h} \frac{\partial h_{\infty}}{\partial z} dz$$
 (3-8)

Incorporation of identities (3-6), (3-7), and (3-8) into equation (3-1) gives

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(3-4)

$$k_{z} \frac{\partial^{2} T}{\partial z^{2}} dA_{d} dz - \dot{m}_{r} C_{pr} \frac{\partial T}{\partial \phi} d\phi - C_{pw} \frac{\partial (m_{w} T)}{\partial \phi} d\phi - (3-9)$$

 $\dot{m}_{a} \frac{\partial h_{m}}{\partial z} dz = 0 .$

In a similar fashion equation (2-10) may be expressed in the form

$$\hat{n}_{a} \frac{\partial h_{\infty}}{\partial z} dz = \frac{f}{C_{pm}} (h_{\infty} - h_{s}) dAv , \qquad (3-10)$$

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(3-12)

where the surface area density is again taken into account to obtain the available area for convection.

The boundary conditions for enthalpy remain the same as those given by equations (2-8) and (2-9).

To transform equations (3-9) and (3-10) into finite difference forms, the differentials may be substituted by finite increments by referring to the grid system developed earlier.

The amounts of air entering the hot and cold sides of an element of matrix are, respectively, $\frac{m_{ah}}{h}$ and $\frac{m_{ac}}{h}$

The convection area for each element equals the total area for convection of one side of the matrix divided by the total number of elements on that size, viz.,

$$\Delta Av_{h} = \frac{Av_{h}}{N_{r} N_{h}} \quad \text{and} \quad \Delta Av_{c} = \frac{Av_{c}}{N_{r} N_{c}} ; \quad (3-11)$$

$$\Delta A_{dh} = \frac{A_{dh}}{N_{c}} \quad \text{and} \quad \Delta A_{dc} = \frac{A_{dc}}{N_{c}} . \quad (3-12)$$

The length of each element is the total length of the matrix divided by the number of rows of the matrix, or

$$z = L/N_{1}$$
, (3-13)

and it is the same for each side of the matrix.

Figure 6 represents an enlarged view of a portion of the overall grid system, and indicates the assigned values of temperatures and enthalpies.

It is recognized that

$$\frac{d^{2}T}{dz^{2}} = \frac{d}{dz} \begin{bmatrix} \lim_{\Delta z \to 0} \frac{T(z) - T(z - \Delta z)}{\Delta z} \end{bmatrix} =$$

$$\lim_{\Delta z \to 0} \frac{1}{\Delta z} \begin{bmatrix} \frac{T(z + \Delta z) - T(z)}{\Delta z} - \frac{T(z) - T(z - \Delta z)}{\Delta z} \end{bmatrix}.$$
(3-14)

When equation (3-14) is applied to the dashed element of Figure 7 one obtains

$$\frac{\mathbf{I}^{2} \mathbf{T}}{\mathbf{I} \mathbf{z}^{2}} = \lim_{\Delta \mathbf{z} \to 0} \left[\left(\frac{\mathbf{T}(\mathbf{i}-\mathbf{1},\mathbf{j}) + \mathbf{T}(\mathbf{i}-\mathbf{1},\mathbf{j}+\mathbf{1})}{2} - \frac{\mathbf{T}(\mathbf{i},\mathbf{j}) + \mathbf{T}(\mathbf{i},\mathbf{j}+\mathbf{1})}{2} \right) - (3-15) \right]$$

$$\left(\frac{\mathrm{T}(\mathrm{i},\mathrm{j})+\mathrm{T}(\mathrm{i},\mathrm{j}+1)}{2}-\frac{\mathrm{T}(\mathrm{i}+1,\mathrm{j})+\mathrm{T}(\mathrm{i}+1,\mathrm{j}+1)}{2}\right)\right]\frac{1}{\Delta z^2}$$

By similar logic the circumferential temperature gradient may be expressed by

$$\frac{dT}{d\phi} = \lim_{\Delta\phi\to 0} \frac{T(i,j+1) - T(i,j)}{\Delta\phi} .$$
(3-16)

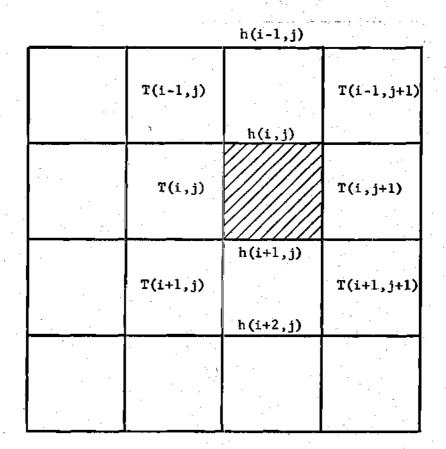


Figure 6. Enlarged View of a Segment of Developed Matrix

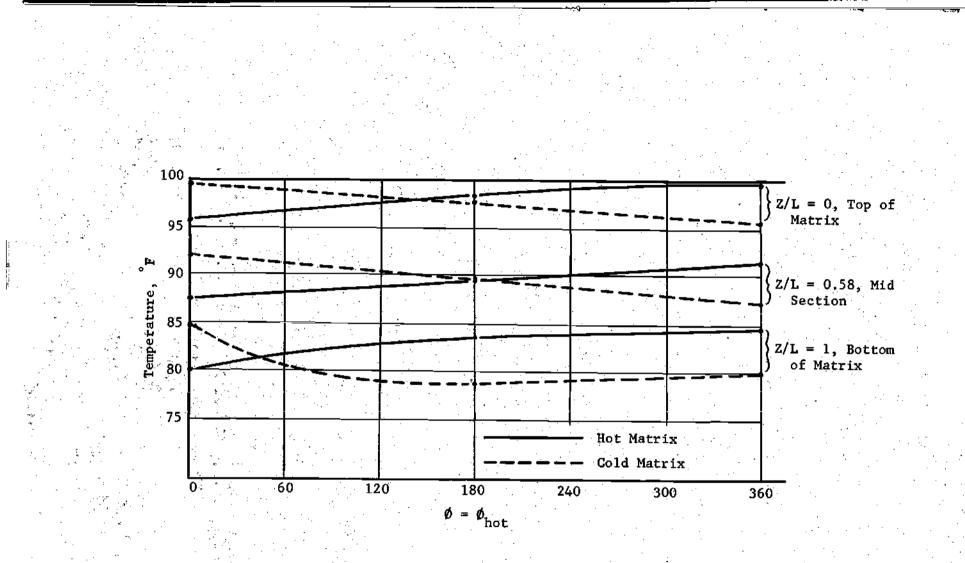


Figure 7. Angular Displacement Along Developed Matrix, ϕ , Degrees

The circumferential gradient of the quantity (m_W^T) may also be expressed, with a few algebraic manipulations, as

$$\frac{d(\dot{m}_{W}T)}{d\phi} = \lim_{\Delta\phi\to 0} \frac{\dot{m}_{W}(\phi + \Delta\phi) T(\phi + \Delta\phi) - \dot{m}_{W}(\phi) T(\phi)}{\Delta\phi} = (3-17)$$

$$\lim_{\Delta\phi\to 0} \left[\frac{\dot{m}_{W}(i, j+1) T(i, j+1) - \dot{m}_{W}(i, j) T(i, j)}{\Delta\phi} \right].$$

Likewise,

$$\frac{dh_{\infty}}{dz} = \lim_{\Delta z \to 0} \left(\frac{h_{\infty}(i+1,j) - h_{\infty}(i,j)}{\Delta z} \right) .$$
(3-18)

Incorporation of equations (3-12) through (3-18) into equations (3-9) and (3-10) yields

$$\frac{N_{r}}{N_{x}} \frac{k}{2} \frac{A_{d}}{L} \left[(T(i-1,j) + T(i-1,j+1) - 2(T(i,j) + T(i,j+1)) + (3-19) \right]$$

$$(T(i+1,j) + T(i+1,j+1) - \frac{m_{r}}{N_{r}} \frac{C_{pr}}{N_{r}} \left[T(i,j+1) - T(i,j) - T(i,j) \right]$$

$$C_{pw} \left[\frac{m_{w}}{m_{w}}(i,j+1) - T(i,j+1) - \frac{m_{w}}{m_{w}}(i,j) - T(i,j) - \frac{m_{w}}{N_{x}} \left[\frac{h_{w}}{m_{w}}(i+1,j) - h_{w}(i,j) \right] \right]$$

A further simplification is

$$N_{h} = N_{c} = N_{r} = N$$

(3-20)

Incorporation of equation (3-20) and a rearrangement yields

$$\frac{\dot{m}_{r} C_{pr}}{N_{r}} \left[T(i, j+1) - T(i, j) \right] + \left[\dot{m}_{w}(i, j+1) C_{pw} T(i, j+1) - (3-21) \right]$$

$$\dot{m}_{w}(i, j) C_{pw} T(i, j) = \frac{\dot{m}_{a}}{N_{x}} \left[h_{\omega}(i, j) - h_{\omega}(i+1, j) \right] + \frac{kA_{d}}{2L} \times \left[(T(i-1, j) + T(i-1, j+1)) - 2(T(i, j) + T(i, j+1)) + (T(i+1, j) + T(i+1, j+1)) \right] = 0.$$

Similar manipulations of equation (3-10) yield

$$\dot{m}_{a}\left[h_{\infty}(i+1,j) - h_{\infty}(i,j)\right] = \frac{fAv}{c_{pm}N_{x}}\left[\frac{h_{\infty}(i,j) + h_{\infty}(i+1,j)}{2} - (3-22)\right]$$

h_{savg}.

Equations (3-21) and (3-22) can be readily applied to any specific element in the matrix. If all the entering conditions of an element are assumed, the problem resolves to determining the unknown exit conditions. In order for this system of equations to be compatible and determinate, four more equations are needed. These are:

a) Conservation of mass equations

$$\frac{m_{a}}{N_{x}} \left[W(i,j) - W(i+1,j) \right] = m_{w}(i,j+1) - m_{w}(i,j) , \qquad (3-23)$$

$$_{a} \left[W(\mathbf{i},\mathbf{j}) - W(\mathbf{i}+1,\mathbf{j}) \right] = \frac{fAv}{C_{pm} N_{r}} \left[\frac{W(\mathbf{i},\mathbf{j}) + W(\mathbf{i}+1,\mathbf{j})}{2} - W_{savg} \right] .$$
 (3-24)

b) Psychrometric equations

$$h = h(T,P)$$
, (3-25)
 $W = W(T,P)$. (3-26)

and

and

m

Equations (3-21) through (3-26) may be applied to the cold side or to the hot side of the matrix by appropriate replacement of subscript x by c or h, respectively.

This set of equations is further simplified (see Appendix I) to yield

$$h_{\infty}(i+1,j) = h_{\infty}(i,j) - HC81 [W(i,j) - W_{savg}(i,j)] C_{pw} T(i,j) (3-27)$$

 $[HC82 m_w(i,j) C_{pw} + HC83][T(i,j+1) - T(i,j)] -$

HC84 [(T(i-1,j) + T(i=1,j+1)) - 2(T(i,j) + T(i,j+1) + (T(i+1,j) + (T(i+1,j)) + (T(i+1,j))) + (

T(i+1, j+1))],

and

$$h_{\infty}(i+1,j) = h_{\infty}(i,j) + HC81 [h_{savg}(i,j) - h(i,j)].$$
 (3-28)

In equations (3-27) and (3-28)

$$HC81 = \frac{2B}{N + B}, \qquad (3-29)$$

$$HC83 = \frac{1 \text{ pr}}{\dot{m}_{a}}, \qquad (3-31)$$

$$B = \frac{fAv}{2 \text{ m}_{a} \text{ C}_{pm}}, \qquad (3-32)$$

$$HC84 = \frac{k \text{ A}_{d}}{2 \text{ m}_{a} \text{L}}. \qquad (3-33)$$

and

A Note on Uniqueness and Stability of the Solution

 $HC82 = \frac{N}{m_a}$

Equations (3-27) and (3-28) have two unknowns, $h_{\infty}(i+1,j)$ and T(i+1,j). It may be seen that, in equation (3-27), $h_{\infty}(i+1,j)$ diminishes when T(i,j+1) increases, and, in equation (3-28), $h_{\infty}(i+1,j)$ increases when T(i,j+1) increases. This implies that there is only one solution which is unique, provided that the factor $\left(1 - \frac{N-B}{N+B}\right)$ in equation (A-7) in Appendix I is greater than one. A negative value of this factor corresponds to an unrealistic physical conclusion of both the hot air and the matrix losing heat simultaneously. Therefore, a greater value than one of the factor is a valid criterion for a stable solution. In other words, the number of elements into which the matrix is subdivided has to be such that $\left(1 - \frac{N-B}{N+B}\right)$ is greater than one, or N > B_h. A similar argument for the cold region leads to the condition N > B_a.

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(3-30)

CHAPTER IV

PROCEDURAL CONSIDERATIONS FOR NUMERICAL SOLUTIONS

Selection of Numerical Values of Input Parameters

Prior to providing an explanation of the numerical procedure adopted to solve the equations developed in the preceding chapter, it is appropriate to discuss the choice of numerical values assigned to several parameters in this investigation.

The previously stated purpose of this study was to demonstrate the validity of the present numerical technique to analyze the RHE rather than obtain design data for a specific application. It may be pointed out that the adaptability of the present technique is not restricted to a unique set of numerical values assigned to various design parameters.

The states of the entering hot and cold streams depend on the specific applications of the RHE. The numerical values selected for the states of entering hot and cold air (hot air at 100 \degree F, 60 percent relative humidity; cold air at 80 \degree F, 50 percent relative humidity), to illustrate the procedure of integration, are within the general range normally found in such applications of the RHE as an air conditioning system and a cloth drying machine. An immediate consequence of the preceding choice of states of entering streams is the specification of their respective enthalpies, specific heats, and specific humidities.

The practical convenience of having approximately the same mass

flow rate for dry streams of hot and cold air, together with the ease of comparing trends of the results of this investigation with those reported by Kays and London (7), dictate the choice of the numerical value for the ratio of minimum to maximum heat capacity rates, C_{\min}/C_{\max} , as 0.7. To ascertain its influence, a value of 0.9 for this parameter was also examined. The mass flow rates for each of the fluid streams consistent with these considerations is approximately 1500 lbm/hr.

Two values for the ratio of the heat capacity rate of the matrix to the minimum heat capacity rate of air, C_r/C_{min} , are considered sufficient to examine the influence of this parameter on the effectiveness of the RHE. These values were selected as two and five, following Kays and London (7).

Use of Psychrometric Equations

Before the numerical technique can be applied to the set of equations derived in Chapter III, it is necessary to have the properties of interest defined by equations of such forms as represented by equations (3-25) and (3-26). Information from steam tables (11) and psychrometric tables (12) can be conveniently used for these formulations. The water vapor properties of interest are the enthalpy and the specific volume as functions of pressure and temperature. Also needed is the saturated water vapor properties in a convenient form to adapt into the scheme of computation. It may be noted that the total pressure used in a psychrometric equation is standard atmospheric pressure, or 14.696 $1bf/in^2$.

It has been assumed that air behaves as an ideal gas yielding the

enthalpy

h = h(T).

Although water vapor deviates from ideal gas behavior, fortunately, the air-water vapor mixture can be still treated as an ideal mixture, with the provision that both air and water vapor occupy the same total volume at the same time. Nevertheless, for the purpose of calculating the dew point temperature, it is reasonable to assume an ideal gas behavior for water vapor. The specific humidity is then given by

$$V = \frac{v_a}{v_w} = C_1 \frac{T}{(P - P_w) v_w} = f(T, P_w)$$
 (4-2)

In equation (4-2), C_1 is a constant with its numerical value determined by the choice of units, B is the total pressure of the mixture, and P_w is the partial pressure of the water vapor. Equation (4-2) indicates that the specific humidity is a function of temperature and partial pressure of water only.

The enthalpy of the air-water vapor mixture can be expressed as

 $h = h_a + w h_w = g(T, P_w)$, (4-3)

which, again, indicates that the enthalpy of the mixture is a function of temperature and partial pressure of water vapor. The datum temperature for enthalpies can be taken at the standard value of $32^{\circ}F$. Finally, the dew point temperature of the mixture, as a function of temperature and specific humidity, can be determined by knowing the correspondingly par-

tial pressure of water vapor and by employing the relation

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(4-1)

where C_2 is a constant with a value dependent on the particular choice of units. The dew point temperature can be evaluated by identifying the saturation temperature from the steam tables corresponding to the known partial pressure of water vapor.

Numerical Integration Procedure

 $w = C_2 \frac{F_W}{(P - P_W)},$

In order to proceed with the solution of the problem, it is necessary to determine the exit thermodynamic properties of the air streams through the matrix. The first step is to obtain the exit conditions for a single matrix element. A step procedure to obtain this may be outlined as

a) Assume the matrix temperature at the exit of an element T(i, j+1) = T(i, j).

c)

b) Evaluate the average surface temperature of the given element. The average humidity ratio at the surface of the element is

evaluated, according to equation (4-2) if the average surface temperature is below the dew point temperature of the entering air; if that is not the case, the average humidity ratio is identical to the humidity ratio of the entering air. In a similar manner, the average enthalpy of the air at the surface of the element is evaluated.

d) Substitution of the above values in equations (3-27) and (3-28) yields two values for the enthalpy, $h_{\omega} = h_{\omega}(i+1,j)$.

e) If the two enthalpies are not identical, within .04 percent, a new T(i, j+1) is substituted and the process is repeated until both equa-

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(4-4)

tions yield the same value of enthalpy, within the acceptable error.

Having obtained T(i, j+1) and $h_{\infty}(i+1, j)$, the rest of the exit properties of the element may be calculated with the help of equations (3-23) through (3-26).

The next step is to assume a distribution of temperatures at the left edge of the hot side of the matrix; i.e., an assignment of the temperatures T(1,1), T(2,1) $T(N_r,1)$ (see Figure 5). Since the entrance conditions of the hot air for the element (1,1) have been obtained, assuming the temperature T(1,1) above and the initial entering condition of the hot air in the matrix, the procedure to obtain the exit conditions for the element (1,1) is similar to the one outlined in the preceding. Obviously, the exit condition of enthalpy for element (1,1) becomes the entering condition of the hot air for element (2,1). The procedure can be repeated to obtain the distribution of enthalpies along the first column of elements and the distribution of matrix temperatures for the right side of the same column of elements. In a similar fashion, the remaining columns of the hot side of the matrix are solved.

To solve for the cold sector side of the matrix, one needs only to adopt the procedure used for the hot sector side, noting that the temperature distribution on the left cold side of the matrix must equal the temperature distribution on the right hot side of the matrix, and that the first element to be calculated is the one at the left bottom of the cold side, where the entrance conditions are known. When the cold side of the matrix is solved, the reversal condition must be satisfied within 0.01 percent. This condition demands that

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In the event that equation (4-5) is not satisfied, the temperature distribution originally assigned to the left edge of the hot side of the matrix is replaced by the calculated temperature distribution at the right edge of the cold side of the matrix. The procedure is repeated to eventually satisfy the reversal condition.

CHAPTER V

PRESENTATION AND DISCUSSION OF RESULTS

A Note on the Presentation of Results

To help attain a physical understanding of the results, all steady state property distributions are presented in matrix form, the latter being identical to the developed matrix of the rotary regenerator. Distributions for the hot and cold sides are represented in separate tables, each of which represents approximately one half of the developed matrix. Among the property distributions presented in the results are the temperature of the matrix, the enthalpy of air, the humidity ratio of air, and the mass rate of condensed water, as well as the states of entering and leaving hot and cold streams. Specific cases are examined for two air flow heat capacity rates, two matrix heat capacity rates, and two heat conductance ratios, the results being examined with respect to the number of heat transfer units of the rotary heat exchanger and heat exchanger effectiveness. A particular set of numerical information to illustrate the results of this study is shown in Table 1.

Property Distributions within the Matrix

Tables 2 and 3, respectively, show the distribution of enthalpies of the hot and cold air streams within the matrix. The top row of Table 2 and the bottom row of Table 3 represent the constant enthalpies of hot and cold air entering the matrix. As the hot stream passed down the

Table 1. Illustrative Set of Numerical Values

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Input Parameters and Properties

Number of rows into which the matrix is divided: 10

40

 $C_{\min}/C_{\max} = 0.7$

 $C_r/C_{min} = 2$ NTUo = 8

Relative humidity of hot entering air = 60%Temperature of hot entering air = $100^{\circ}F$ Enthalpy of hot entering air = 44.07 Btu/1bm Relative humidity of cold entering air = 50%Temperature of cold entering air = $80^{\circ}F$ Enthalpy of cold entering air = 23.48 Btu/1bm

Output_Properties

Enthalpy of hot leaving air = 39.76 Btu/1bm Enthalpy of cold leaving air = 28.06 Btu/1bm

Computer Time = 4.51 minutes

Table 2. Enthalpy of Air in the Hot Part of the Developed Matrix

· ·	· .		•						
44.07	44.07	44.07	44.07	44.07	44.07	44.07	44.07	44.07	44+07
43.44	43.56	43.72	43.77	43.79	43.85	43.94	43.97	43.97	43+97
43.08	43.13	43.31	43.36	43.46	43.53	43+60	43.72	43.76	43.78
42.65	42.78	42.85	43.00	43.10	43.19	43.27	43.37	43.44	43.48
42.21	42.39	42.43	42.60	42.67	42.82	42-91	43.02	43.10	43•16
41.79	41+90	42.00	42.19	42.25	42.36	42.53	42.59	42.73	42.82
41.33	41+44	41.56	41.70	41.81	41.91	42.05	42.15	42.26	42+38
40.77	40.89	41+07	41.15	41.32	41.45	41.57	41.69	41+81	41.92
40.08	40+61	40.39	40.58	40.74	40.88	41.01	41.14	. 41.32 .	41+39
38.09	39.61	. 40+17	40.42	40.62	40.79	40.94	40.44	40.62	40+81
36,41	38.13	39.32	39.98	40.36	40.61	40+80	40.68	40.62	40•74
	-	100 A. A. A.							

- NOTES: a) Hot air enters at top left.
 - b) Columns indicate angular positions along the matrix circumference at an interval of 20° each.
 - c) Rows indicate axial positions along the matrix length.

Table 3. Enthalpy of Air in the Cold Part of the Developed Matrix

30.11941 28.09440 28.03082 27.98095 27.89512 27.88261 27.77042 27.65736 27.64417 27.50429 29.92846 27.88260 27.77961 27.75253 27.60788 27.54742 27.45385 27.33955 27.30299 27.14836 29.59091 27.58000 27.47357 27.39103 27.25121 27.21777 27.08543 27.01155 26.93201 26.77775 29.25540 27.22675 27.08247 27.03956 26.90486 26.83170 26.74611 26.58013 26.50386 26.40630 28.89594 26.85779 26.72359 26.65156 26.56113 26.39365 26.30797 26.19351 26.07412 25.94614 28.51390 26.47197 26.38500 26.21412 26.11530 26.00741 25.85912 25.70920 25.59135 25.44549 28.13339 26.07680 25.93262 25.81934 25.64861 25.54083 25.32017 25.17512 25.02545 24.83532 27.98904 25.70717 25.56681 25.32165 25.07463 24.89963 24.68665 24.48729 24.34230 24.24949 28.71479 25.25908 24.84513 24.47966 24.24273 24.01622 23.90554 23.73759 23.66657 23.63707 33.30295 22.61549 22.74968 22.92937 22.99918 23.10912 23.13496 23.21864 23.26458 23.29309 23.48267 23.

NOTES: a) Cold air enters at lower left.

- b) Columns indicate angular positions along the matrix circumference at an interval of 20° each.
- c) Rows indicate axial positions along the matrix length.

passages in the matrix it delivers energy to the metal matrix elements, a process indicated by a decrease in enthalpy. In a similar fashion the cold stream entering at a lower level of enthalpy at the bottom of the cold part of the matrix primarily receives energy from the hotter metal matrix, this process resulting in an increase of enthalpy of the cold stream as it moves up through the matrix. The bottom row of Table 2 and the top row of Table 3, respectively, represent the enthalpies of the hot and cold streams leaving the matrix. Average values of these exit enthalpies are shown in Table 1.

Tables 4 and 5, respectively, show the distribution of temperature within the hot and cold sides of the matrix. The same information, for selected rows of the entire matrix, is shown graphically in Figure 7. The predominant variation of temperature is seen to be along the rows in these two tables, since these correspond to the rotation of the matrix about its axis. The last column of Table 4 also represents the first column of Table 5, since it represents the state of the matrix column leaving the hot domain and entering the cold domain. The trend, previously explained in connection with the discussion on enthalpy distribution, is visible in the temperature distribution as well; viz, a rise in matrix temperature due to "energy received from the hot air stream, with a subsequent decrease in matrix temperature due to energy transferred to the cold stream. The last column of Table 5 is the input information for iteration of the property distributions. Ideally, this column should be identical to the first column of Table 4, the actual difference being the allowed truncation error in satisfying the so-called reversal condition.

Table 4. Temperature of the Hot Part of the Developed Matrix

								•	· ·		
	96.52	97.20	97+74	98+12	98.44	98,75	98-98	99+12	99+23	99+33	99.43
	95.22	95.62	96.08	96+52	96.95	97.31	97+66	98+02	98+28	98.50	98+71
	93.63	94.09	94.47	94.97	95.35	95.73	96.10	96.47	96.84	97+19	97+50
	91.98	92.45	92+87	93.32	93.75	94+21	94+61	94.99	95.37	95+73_	96.07
	90.15	90.60	91.12	91.58	92.02	92.48	92+97	93,39	93+85	94.24	94+60
	88.16	88+66	89+15	89+63	90.16	90+63	91+12	91.63	. 92+10 .	92+60	93+08
	85.84	86.44	87+04	87.56	88.14	88.66	89+15	89+65	90+14	90+63	91+12
•	83.20	83.93	84.26	84.98	85.60	86+23	86+84	87.44	88.03	88+56	89+12
	80.55	82+58	83+62	83.86	84.05	. 64.21	84+35	84.48	85.23	85+97	86-59
	79.25	80.95	62.46	. 83.32	83.76	84.02	64+21	84+36	84.14	84+18	84+30
			1. Contract 1. Con								

- NOTES: a) Hot air enters at top left.
 - b) Columns indicate angular positions along the matrix circumference at an interval of 20° each.
 - c) Rows indicate axial positions along the matrix length.

					_	-	
· · ·	· .						
				-			•
99.43 99.24 99.02	98.76 98.53	98,23	97-69	97.57	97.24	96+89	96.53
· · · · · · · · · · · · · · · · · · ·	· · · · · · · · · · · · · · · · · · ·		_			_	
98.71 98.36 98.06	97.74 97.37	97.02	96+68	96+31	95+97	95+59	95+22

Table 5. Temperature of the Cold Part of the Developed Matrix

	70+/*	70.30	70+00	91+14	91.31	9/102	A0+09	90+31	30+31	42+24	95+22	
• .	97.50	97.16	96.80	96.40	96+04	95.70	95+30	94+96	94+52	94+08	93+70	
	96+07	. 95.70	95+32	94.96	94.57	94.22	93.77	93.32	92+93	92+49	92.02	
	94.60	94.21	93.82	93.47	93.03	92.57	92.18	91+72	91.22	90+73	90+22	
	93.08	92.69	92.29	91-83	91.43	90.95	90.48	89.93	89+38	88+A1	88+19	
	91.12		90+60	90.22	89.72	89.13	88.48	87.84	87.14	86+45	85+85	•
	89.12	89.86	89+40	88.67	87.81	85.97	86+08	85+29	84+53	83+85	83+23	
	86.59	91.23	88+55	86.43	84.86	63.60	62+68	A1.90	.81+38	80.97	80+62	
	84+30		76+03		77.33	. 77.83	78.21	78.56	78+83	79+05	79+25	
	1											

- NOTES: a) Cold air enters at lower left.
 - b) Columns indicate angular positions along the matrix circumference at an interval of 20° each.
 - c) Rows indicate axial positions along the matrix length.

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The distribution of humidity ratio of the hot and cold air is, respectively, shown in Tables 6 and 7. It is seen from Table 6 that, except for the bottom three rows, the humidity ratio is constant through the hot domain. This implies that condensation of water in the hot domain occurs only within the bottom three rows. The actual amount of condensation can be interpreted in terms of the decrease in humidity ratio of hot air along the three rows. On the other hand, only the first column of Table 7 indicates any change in the humidity ratio. In fact, this occurs only on the bottom element of the first column of Table 7. An account of this behavior might be provided by the carry-over of condensed water from the hot side to the cold side, a subsequent transfer of energy from this water to the relatively cooler entering air, and an eventual evaporation of all the carried water owing to heat absorbed from the matrix. There is no further increase in specific humidity of the first column of the cold side.

A qualitative verification of the explanation regarding carry-over can be sought by examining Tables 8 and 9, which respectively represent the distribution of the rate of condensed water within the hot and cold sides of the matrix. It is clearly seen from Table 8 that condensation does occur within the bottom three rows and also that the rate of condensation increases to a maximum along the rows, followed by a reevaporation of a portion of the condensate, resulting in an apparent display of a reduced condensation. The portion of condensate that does not get reevaporated on the hot side is carried over to the cold side as explained previously. Table 6. Specific Humidity of the Hot Part of the Developed Matrix

e.

· •									· · · ·
.02511	+02511	.02511	.02511	.02511	.02511	.02511	.02511	.02511	.02511
.02511	•02511	.02511	+02511	.02511	.02511	+02511	•025 1 1	.02511	.02511
.02511	.02511	.02511	.02511	+02511	.02511	.02511	.02511	.02511	.02511
.02511	.02511	.02511	+02511	.02511	.02511	.02511	.02511	.02511	.02511
.02511	.02511	.02511	.02511	•02511	+02511	.02511	.02511	.02511	.02511
-02511	.02511	.02511	.02511	+02511	.02511	.02511	.02511	•02511	.02511
+02511	+02511	.02511	.02511	+02511	.02511	.02511	.02511	.02511	.02511
+02511	+02511	.02511	.02511	.02511	+02511	.02511	.02511	.02511	.02511
.02510	.02513	.02511	.02511	.02511	.02511	.02511	.02511	.02511	.02511
.02377	.02478	.02522	.02536	.02548	•02559	.02557	.02511	•02511	.02511
.02260	.02380	.02466	•02512	.02539	•02556	•02568	.02557	.02551	.02556

- NOTES: a) Hot air enters at top left.
 - b) Columns indicate angular positions along the matrix circumference at an interval of 20° each.
 - c) Rows indicate axial positions along the matrix length.

Table 7. Specific Humidity of the Cold Part of the Developed Matrix

	and the second se								
	•01270•01092•01092	+01092	.01092	.01092	.01092	.01092 .	.01092	+01092	
	•01270 •01092 •01092	•01092	+01092	.01092	.01092	-01092	.01092	.01092	
	.01270	+01092	.01092	.01092	•01092	.01092	.01092	.01092	
	.012700109201092	+01092	.01092	.01092	.01092	+01092	.01092 .	+01092	
	•01270 •01092 •01092	•01092	.01092	+01092	.01092	+01092	.01092	.01092	
	•01270	•01092	•01092	.01092	+01092	.01092	+01092 🔿	.01092	
	•01270•01092•01092	•01092	+01092	+D1092	01092 -	.01092	+01092	.01092	
•	.01270 .01092 .01092	•01092	.01092	•01092	.01092	+01092	•01092	.01092	
	•01270		•01092	.01092	.01092	+01092	•01092	.01092	
	.01270 .01092	+01092	•01092	.01092	.01092	•01092 · .	•01092	.01092	
	.01092	. • 61092		+01092.	01092	•01092	+01092	.01092	

- NOTES: a) Cold air enters at lower left.
 - b) Columns indicate angular positions along the matrix circumference at an interval of 20° each.
 - c) Rows indicate axial positions along the matrix length.

Table 8. Mass Rate of Condensed Water on Hot Part of the Developed Matrix

									•	
+00000		. 00000	.00000	.00000	.00000	.00000	.00000	-00000	.00000	+00000
•00000	+00000	.00000	• 00000	.00000	.00000	.00000	-00000	.00000	.00000	.00000
+00000	•00000	+00000 .	•00000	.00000	•00000	.00000	.00000	.00000	•a000 0	.00000
•00000	•00000	+00000	+00000	+00000	00000	.00000	•00000	.00000	.00000	+00000
+00000	+00000	•00000	+00000	•00000	+00000	.00000	.00000	•00000	•0000	+00000
•00000	•00000	.00000	•00000	•000000	.00000	.00000	•00000	.00000	+00000	.00000
•00000	+00000	.00000	•00000	.00000	•00000	+00000	+00000	+00000	.00000	+00000
+00000	•00239	•00000 *	+00000	.00000	+00000	.00000	+00000	.00000	+00000	+00000
•00000	•19879	•25049	+23480	.19701	.14098	+06940	•00000	.00000	.00000	.00000
-00000	•17556	.32300	•40681	•44266	•45721	+46106	+44533	.37684		•25023

NOTES: a) Hot air enters at top left.

- b) Columns indicate angular positions along the matrix circumference at an interval of 20° each.
- c) Rows indicate axial positions along the matrix length.

Table 9. Mass Rate of Condensed Water on Cold Part of the Developed Matrix

:									• •	
+00000	•00000	.00000	.000.00	.00000	.00000	.00000	•00000	.00000	.00000	.00000
•00000	.00000	.00000	+60000	.00000	+00000	•00000	.00000	+00000	•0000n	.00000
•00000	+00000	. •00000 .	.00000	•00000	-00000	.00000	+00000	.00000	.00000	•0000n
+00000	. 00000 .	+00000	•00000	.00000	+00000	.00000	•00000	+00000	+00000	.00000
.00000		00000	+00000	+00000	+00000	.00000	+00000	.00000	.00000	•00000
.00000	+00000	.00000	+00000	.00000	.00000	.00000	•00000	.00000	.00000	.00000
+00000	00000 .	+00000	.00000	.00000	+00000	.00000	•00000	.00000	.00000	.00000
-00000	+00000	.00000	.00000	.00000	+00000	.00000	+00000	+00000	+00000	.00000
+00000	+00000	.00000	.00000	.00000	+00000	.00000	•00000	+00000		.00000
.25023			00000	.00000	,00000	.00000	.00000	.00000	.00000	.00000

NOTES: a) Cold air enters at lower left.

b) Columns indicate angular positions along the matrix circumference at an interval of 20° each.

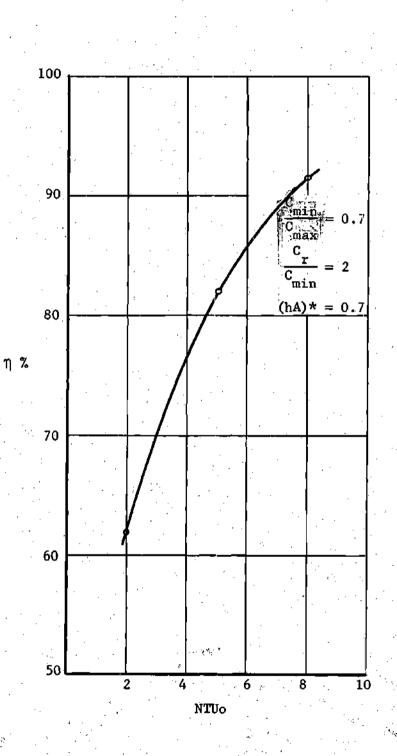
g

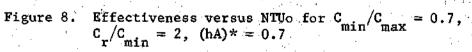
c) Rows indicate axial positions along the matrix length.

Parametric Considerations of the Regenerator Performance

Figures 8 through 11 show the performance curves of the rotary regenerator under the influence of several parameters. Each of these curves was obtained by a careful examination of the data from several sets of eight tables similar to those discussed previously. The rotary regenerator effectiveness and the number of transfer units (NTUo) are the basic variables in these curves, plotted for two values of the ratio of minimum to maximum heat capacity rates, C_{min}/C_{max} , two values of the ratio of rotor heat capacity rate to the minimum heat capacity rate of air, C_r/C_{min} , and two values of the ratio of effective conductance (hA)*.

It is recognized that the number of transfer units, NTU, expresses the "heat transfer size" of a heat exchanger configuration. As such, it is to be expected that the higher value of NTUo should correspond to a greater heat transfer rate, both from the hot air to the matrix as well as from the matrix to the cold air. Such a trend is seen in Figure 8. Further, an examination of Tables 10 and 11, respectively representing the enthalpy distribution for the hot domain for NTUO = 5 and NTUO = 8, and an examination of Tables 12 and 13, respectively representing the cold domain enthalpy distribution for NTU0 = 5 and NTU0 = 8, confirm the expectation regarding the hot and cold domain heat transfer rates. Moreover, an examination of the distribution of condensing rates for NTUo = 5 and NTUo = 8, respectively shown in Tables 14 and 15, indicates a higher condensation rate for the higher NTUo case, an observation consistent with the greater heat transfer rate mentioned earlier. Similar observations regarding the influence of NTUo can be made from Figures 9, 10, and 11.





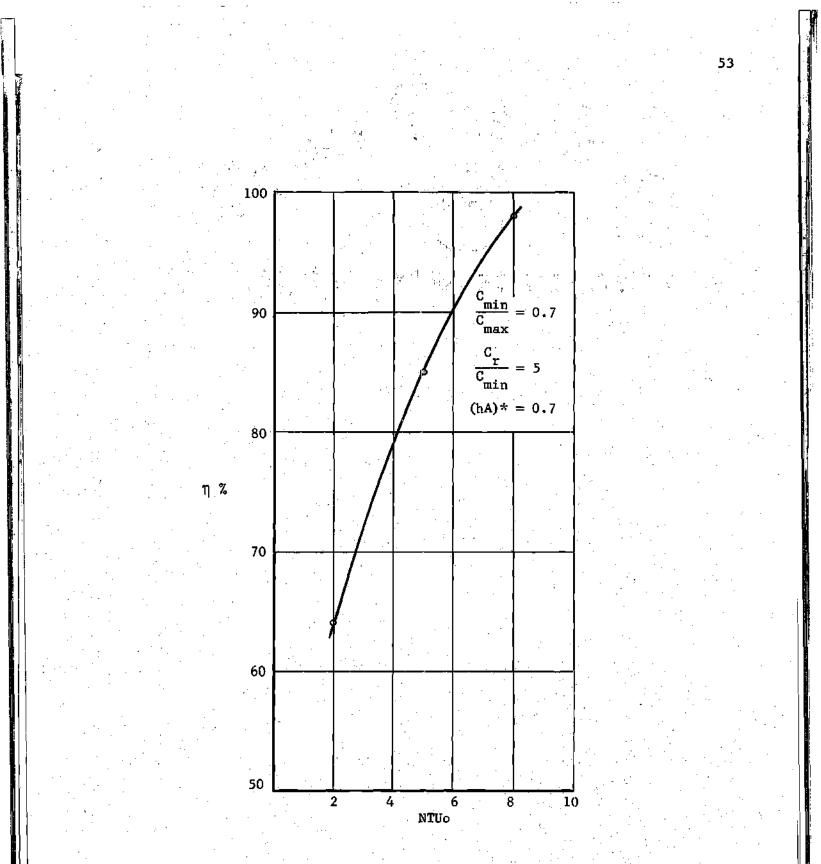
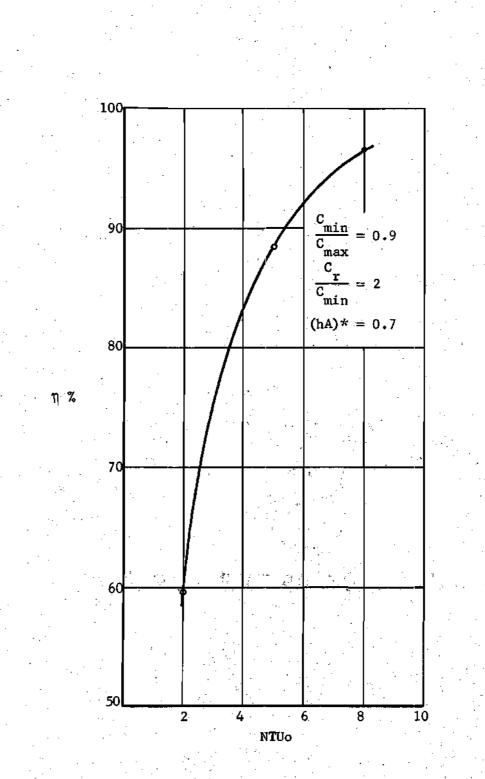
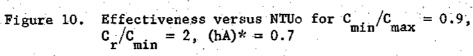
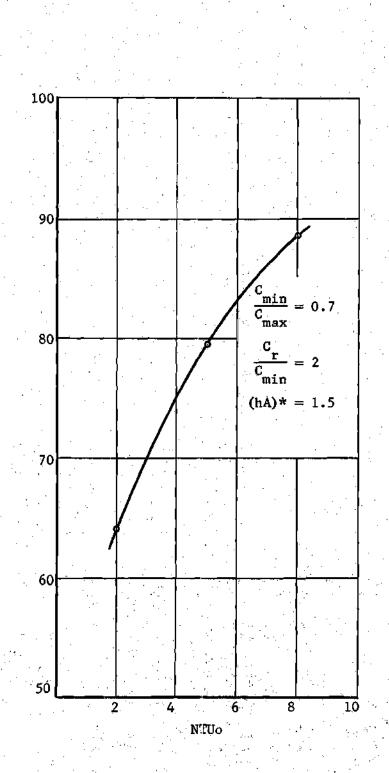


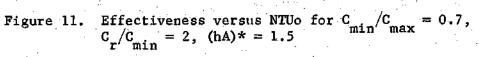
Figure 9. Effectiveness versus NTUo for $C_{min}/C_{max} = 0.7$, $C_r/C_{min} = 5$, (hA)* = 0.7







η %



									· .		
44+07	44.07	44.07	44.07	44.07	44,07	44.07	44.07	44.07	44.07	44.07	
43.60	43.66	43.70	43.74	43.78	43,80	43.84	43,86	43.87	43.91	43.94	
43.18	43.28	45.35	42.41	43+46	43,53	43.55	43.61	43.64	43.69	43.72	
42.80	42.91	43.00	43.07	43.13	43,20	43.27	43.31	43.38	43,45	43.49	
42.43	42.55	42.64	42.72	42.79	42.A6	42.98	43.61	43.09	43,16	43.23	
42,08	42.19	42.29	42.38	42.45	42,53	42.63	42,68	42.77	42,86	42.92	
41+70	41.82	41.93	42.02	42.10	42.18	42.28	42.34	42.43	42.51	42.58	
41.31	41+44	41.52	41.64	41.74	41,83	41.92	42.00	42.08	42.17	42.24	
40.92	41.02 _		41+27	41.34	41,45	41.55	41.64	41.73	41,82	41.90	•.
40.27	40.61	40.86	4 0.60	40.89	41.04	41,13	41.25			41.52	
39.08	39.77	40.28	40.48	40,68	40.48	\$1.01	41.14	41.26	40.98	41+07	
37.76	. 38.71	39.46	34.92	40+28	40.58		40,98		41.01	41+07	

Table 10. Enthalpy of Air in the Hot Part of the Developed Matrix

NOTES: a) Hot air enters at top left. Hot air enters at top left. b) Columns indicate angular positions along the ma-trix circumference at an interval of 20° each. c) Rows indicate axial positions c) Rows indicate axial positions along the matrix length. d) NTUo = 5.

 $\frac{C_{\min}}{C_{\max}} = 0.7$ e) $\frac{C_r}{C_{min}}$ f)

= 2 .

	Table 11.	Enthalpy of Air in	the Hot	Part of the	e Developed Matrix
_					
				·	

				· .						
44.07	44.07	44.07	44+07	44.07	44.07	44.07	44.07	44.07	44+07	44.07
43.55	43.63	43+67	43.74	43.77	43.83	43.86	43+90	43.92		43.97
					43+51	43+60	43+66	43.70	. –	
							43.33	43.43	43+45	43.52
42.33	42.46	.42+58	42.69	42.77	42.84	42+92	43+00		43+13	43.24
41.95	42.09	42.21	42.27	42.39	42.49	42+58	42+66	42.74	42+80	42.89
41+54	41.69	41.78	. 41+90	41.97	42+10	42+21	42+26	42.37	42+46	42.55
41.10	41.23	41+33	41.49	41.55	41.68	41+80	41.90	41.95	42+08	42.10
40.58	40+74	40+84	. 40+97	41.08	41.22	41+34	41.44	41.52	41-68	41.76
0 4•34	40+09	40+50	40.77 :	. 40.96	40.66	40+78				
36.65	37.86	38.87	39+64	40-19	40.45	40.62				
	43.55 43.15 42.73 42.33 41.95 41.54 41.10 40.58 39.34 37.71	43.55 43.63 43.15 43.21 42.73 42.84 42.33 42.46 41.95 42.09 41.54 41.69 41.10 41.23 40.58 40.74 39.34 40.09 37.71 38.95	43.55 43.63 43.67 43.15 43.21 43.31 42.73 42.84 42.95 42.33 42.46 42.58 41.95 42.09 42.21 41.54 41.69 41.78 41.10 41.23 41.33 40.58 40.74 40.84 39.34 40.09 40.50 37.71 38.95 39.77	43.15 43.21 43.31 43.40 42.73 42.84 42.95 43.04 42.33 42.46 42.95 43.04 42.33 42.46 42.58 42.69 41.95 42.09 42.21 42.27 41.54 41.69 41.78 41.90 41.10 41.23 41.33 41.49 40.58 40.74 40.84 40.97 39.34 40.09 40.50 40.77 37.71 38.95 39.77 40.31	43.55 43.63 43.67 43.74 43.77 43.15 43.21 43.31 43.40 43.45 42.73 42.84 42.95 43.04 43.12 42.33 42.46 42.58 42.69 42.77 41.95 42.09 42.21 42.27 42.39 41.54 41.69 41.78 41.90 41.97 41.10 41.23 41.33 41.49 41.55 40.58 40.74 40.84 40.97 41.08 39.34 40.09 40.50 40.77 40.95 37.71 38.95 39.77 40.31 40.67	43.55 43.63 43.67 43.74 43.77 43.83 43.15 43.21 43.31 43.40 43.45 43.51 42.73 42.84 42.95 43.04 43.12 43.16 42.33 42.46 42.58 42.69 42.77 42.84 41.95 42.09 42.21 42.27 42.39 42.49 41.54 41.69 41.78 41.90 41.97 42.10 41.10 41.23 41.33 41.49 41.55 41.68 40.58 40.74 40.84 40.97 41.08 41.22 39.34 40.09 40.50 40.77 40.96 40.66	43.55 43.63 43.67 43.74 43.77 43.83 43.86 43.15 43.21 43.31 43.40 43.45 43.51 43.60 42.73 42.84 42.95 43.04 43.12 43.18 43.26 42.33 42.46 42.58 42.69 42.77 42.84 42.92 41.95 42.09 42.21 42.27 42.39 42.49 42.58 41.54 41.69 41.76 41.90 41.97 42.10 42.21 41.10 41.23 41.33 41.49 41.55 41.68 41.80 40.58 40.74 40.84 40.97 41.08 41.22 41.34 39.34 40.09 40.50 40.77 40.96 40.66 40.78 37.71 38.95 39.77 40.31 40.67 40.67 40.75	43.55 43.63 43.67 43.74 43.77 43.83 43.86 43.90 43.15 43.21 43.31 43.40 43.45 43.51 43.60 43.66 42.73 42.84 42.95 43.04 43.12 43.16 43.26 43.33 42.33 42.46 42.58 42.69 42.77 42.84 42.92 43.00 41.95 42.09 42.21 42.27 42.39 42.49 42.58 42.66 41.54 41.69 41.78 41.90 41.97 42.10 42.21 42.26 41.10 41.23 41.33 41.49 41.55 41.68 41.80 41.90 40.58 40.74 40.84 40.97 41.08 41.22 41.34 41.44 39.34 40.09 40.50 40.77 40.96 40.67 40.75 40.87 35.65 37.86 34.87 20.64 40.47 40.87 40.67 40.75 40.87	43.55 43.63 43.67 43.74 43.77 43.83 43.86 43.90 43.92 43.15 43.21 43.31 43.40 43.45 43.51 43.60 43.66 43.70 42.73 42.84 42.95 43.04 43.12 43.16 43.26 43.33 43.43 42.33 42.46 42.58 42.69 42.77 42.84 42.92 43.00 43.08 41.95 42.09 42.21 42.27 42.39 42.49 42.58 42.66 42.74 41.54 41.69 41.78 41.90 41.97 42.10 42.21 42.26 42.37 41.10 41.23 41.33 41.49 41.55 41.68 41.90 41.95 40.58 40.74 40.84 40.97 41.08 41.22 41.34 41.44 41.52 40.67 40.66 40.78 40.94 40.97 37.71 38.95 39.77 40.31 40.67 40.67 40.75 40.87 40.97 36.65 37.86 34.87 20.64 40.47 40.87 40.97 40.67 40.67 40.75 40.87 40.97	43.55 43.63 43.67 43.74 43.77 43.83 43.86 43.90 43.92 43.95 43.15 43.21 43.31 43.40 43.45 43.51 43.60 43.66 43.70 43.73 42.73 42.84 42.95 43.04 43.12 43.16 43.26 43.33 43.43 43.45 42.33 42.46 42.58 42.69 42.77 42.84 42.92 43.00 43.08 43.13 41.95 42.09 42.21 42.27 42.39 42.49 42.66 42.74 42.80 41.54 41.69 41.78 41.90 41.97 42.10 42.21 42.26 42.37 42.46 41.54 41.69 41.73 41.90 41.97 42.10 42.21 42.26 42.37 42.46 41.54 41.69 41.73 41.90 41.97 42.10 42.21 42.26 42.37 42.46 41.54 41.69 41.78 41.90 41.97 42.10 42.21 42.26 42.37 42.46 40.58 40.74 40.84 40.97 41.68 41.80 41.90 41.95 42.08 40.58 40.74 40.84 40.97 41.08 41.22 41.34 41.44 41.52 41.68 39.34 40.09 40.97 40.96 40.66 40.78 40.94 41.01 41.19 37.71 38.95 39.77

NOTES: a) Hot air enters at top left. Hot air enters at top left. b) Columns indicate angular positions along the matrix circumference at an interval of 20° each. c) Rows indicate axial positions along the matrix length. d) NTUo = 8.

 $\frac{C_{\min}}{C_{\max}} = 0.7 .$ e)

Cmin

f)

Table 12. Enthalpy of Air in the Cold Part of the Developed Matrix

29.45825 27.86223 27.70954 27.60749 27.59201 27.52295 27.43368 27.40775 27.32129 27.22036 27.11858 29.18750 27.58033 27.47554 27.37006 27.32560 27.26581 27.14886 27.06038 27.03686 26.91804 26.81646 28.90315 27.28077 27.10470 27.00533 27.02324 26.94499 26.92769 26.77459 26.69361 26.59156 26.51268 28.59305 26.96096 26.84051 26.76352 26.72457 26.60011 26.50301 26.45880 26.34737 26.28732 26.16237 28.25006 26.62528 26.50947 26.44834 26.36159 26.26391 26.20673 26.08468 26.02898 25.88653 25.82319 27.89745 26.29531 26.18672 20.14549 26.02368 25.96477 25.81399 25.74557 25.63904 25.52634 25.41951 27.60787 25.99359 25.87276 25.77500 25.68906 25.56475 25.44117 25.33888 25.24558 25.13036 24.95646 27.39280 25.68124 25.54342 25.47196 25.25934 25.14908 25.92867 24.88877 24.78294 24.65372 24.54195 27.45120 25.40597 25.2073 25.04272 24.81219 24.66945 24.52936 24.36809 24.30284 24.15949 24.06117 28.49370 24.85164 24.55596 24.33849 24.15263 23.98628 23.84920 23.80877 23.73311 23.62923 23.59491 32.68224 22.84928 22.91016 23.01166 23.07344 23.17457 23.24777 23.27228 23.29504 23.30728 23.37800 23.40267 23.48267 23.4

NOTES: a) Cold air enters at lower left. b) Columns indicate angular positions along the matrix circumference at an interval of 20° each. c) Rows indicate axial positions along the matrix length. d) NTUO = 0.7.

άO.

<u>min</u> Cmax

min

Table 13. Enthalpy of Air in the Cold Part of the Developed Matrix

30.45182 28.11539 28.03641 27.95771 27.87398 27.79877 27.73940 27.65854 27.61307 27.54977 27.42874 30.20107 27.83480 27.76524 27.68866 27.59093 27.55092 27.48971 27.36777 27.30365 27.21263 27.09774 29.68108 27.52278 27.52500 27.39582 27.29329 27.26770 27.13625 27.05463 26.96916 26.83408 26.80135 29.55663 27.20852 27.18998 27.05239 27.01559 26.86822 26.82724 26.67144 26.62865 26.47755 26.43355 29.28877 26.88091 26.84182 26.69012 26.64802 26.49643 26.44832 26.29494 26.24841 26.10553 26.05517 28.68581 26.52401 26.46705 26.31705 26.27287 26.11565 26.06696 25.91764 25.85878 25.76412 25.64125 28.50482 26.11310 26.08371 25.94612 25.89448 25.73226 25.67064 25.56466 25.42888 25.29278 25.13601 28.09480 25.78874 25.70064 25.60915 25.47011 25.36154 25.24334 25.06767 24.91656 24.78236 24.56542 27.70017 25.51615 25.41119 25.32349 25.05936 24.84907 24.63279 24.37675 24.22173 24.12121 23.94906 28.07446 25.68039 25.01629 24.54704 24.15609 23.87365 23.70330 23.59164 23.57637 23.51112 23.41061 34.69885 22.29870 22.52964 22.72435 22.86099 23.00829 23.08641 23.20763 23.26562 23.22726 23.48267

e)

f)

min

min

Cmax

= 0.7

NOTES: a) Cold air enters at lower left. b) Columns indicate angular positions along the matrix circumference at an interval of 20° each. c) Rows indicate axial positions along the matrix length. d) NTUo = 8.

Ś

Table 14. Mass Rate of Condensed Water on the Hot Part of the Developed Matrix

		·	•				·.	· · ·		1.1	
.00000	.00000	.00000	.00000	.00000	.00000	.00000	.00000	.00000	.00000	.00000	.00000
.00000	.00000	. 00000	.00000	+00060	.000n0	.00000	.00000	.00000	.00000	.00000	+00000
.00000	.00000			.00000	.00000	.00000	.00000	+00000	•00000	.00000	-00000
+00000	+00000	.00000	.00000	.00000	.00000	.00000	.00000	.00000	•00000	.00000	+00000
.00000	.00000	.00000	.00000	.00000	.00000	.00000	.00000	.00000	•00000	.00000	•00000
•00000	.00000	.00000	.00000	+00000	+00000	•00000	+00000	•00000	+00000	.00000	+00000
.00000	.00000	.00000	.00000	.00000	+00000	+00000	•00000	+00000	+00000	.00000	+00000
+00000	.00000	.00000	.00000	+00000	.00000	.00000	+00000	+00000	•00000	.00000	•00000
+00000	.02285		.00000	. 00000	+00000	.00000	.00000	+00000	•00000	.00000	•00000
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NOTES: a) Hot air enters at top left. b) Columns indicate angular positions along the matrix circumference at an interval of 20° each. c) Rows indicate axial positions along the matrix length. d) NTUo = 5.

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e) $\frac{C_{\min}}{C_{\max}} = 0.7$. f) $\frac{C_{\max}}{C_{\min}} = 2$. Table 15. Mass Rate of Condensed Water on the Hot Part of the Developed Matrix

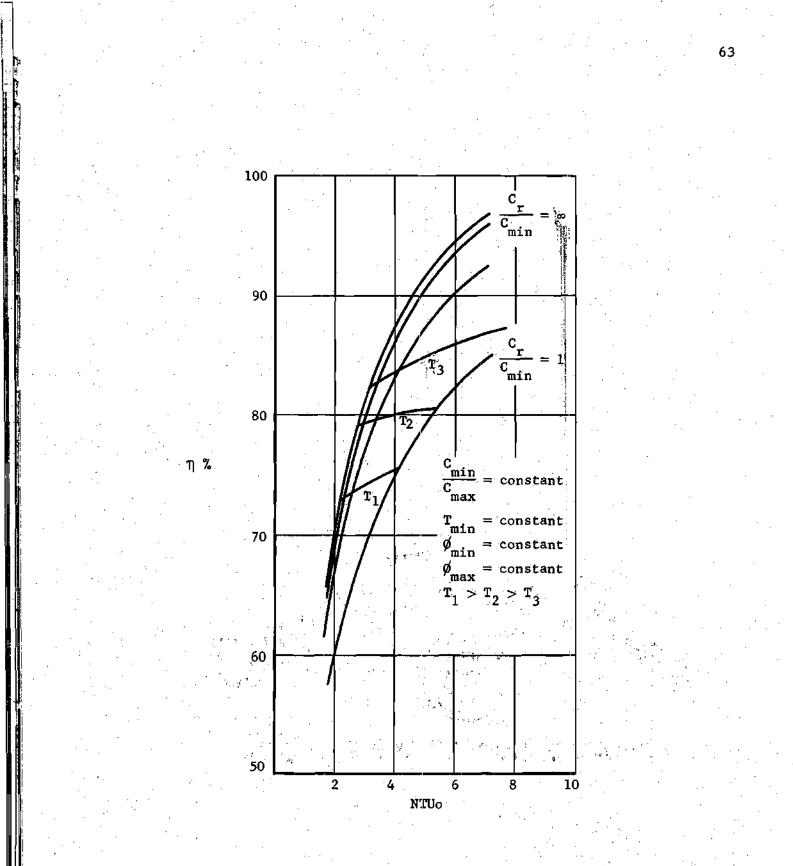
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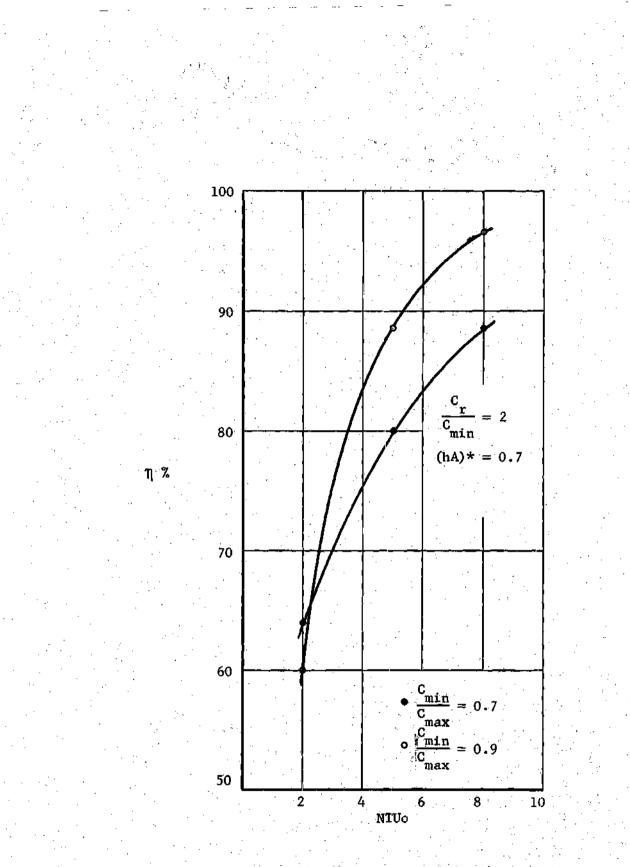
NOTES: a) Hot air enters at top left. b) Columns indicate angular positions along the matrix circumference at an interval of 20° each. c) Rows indicate axial positions along the matrix length. d) NTUo = 8.
e) Cmin = 0.7.
f) Cr = 2.

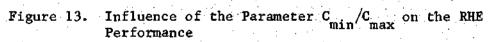
A very interesting observation regarding the influence of condensation on the rotary heat exchanger effectiveness can be made on the basis of the preceding discussion on the role played by NTUo in the heat transfer process. It appears that the start of cold air humidification, which might prove to be a critical design condition in applying the RHE to such systems where it is to be avoided, can be identified as a point on each of the NTUo - Effectiveness curves for a given set of such parameters as C_{min}/C_{max} , C_r/C_{min} , and (hA)*. In other words, although the effectiveness shows an apparent increase due to the cold air humidification, the need to avoid the latter puts an upper limit on the former. Figure 12 shows qualitatively how such design information might appear on the Effectiveness versus NTUo curve for prescribed states of entering air streams. That such information was not obtained in this investigation due to its previously stated limited scope, does not preclude a further exploration of the numerical model developed herein just to obtain such data.

Figure 13 shows the influence of the parameter C_{min}/C_{max} on the RHE effectiveness. It is seen that, for the case of a higher value of the heat capacity rate ratio, the effectiveness is higher, both for NTUo = 5 and for NTUO = 8. A possible explanation for such a trend can be provided by first observing that a higher value of C_{min}/C_{max} amounts to a higher flow rate of cold air in this investigation. As a result, the average temperature distribution within the matrix is at a lower level than that for the case of $C_{min}/C_{max} = 0.7$. The consequence of this is that, for the higher average temperature distribution, all the condensate









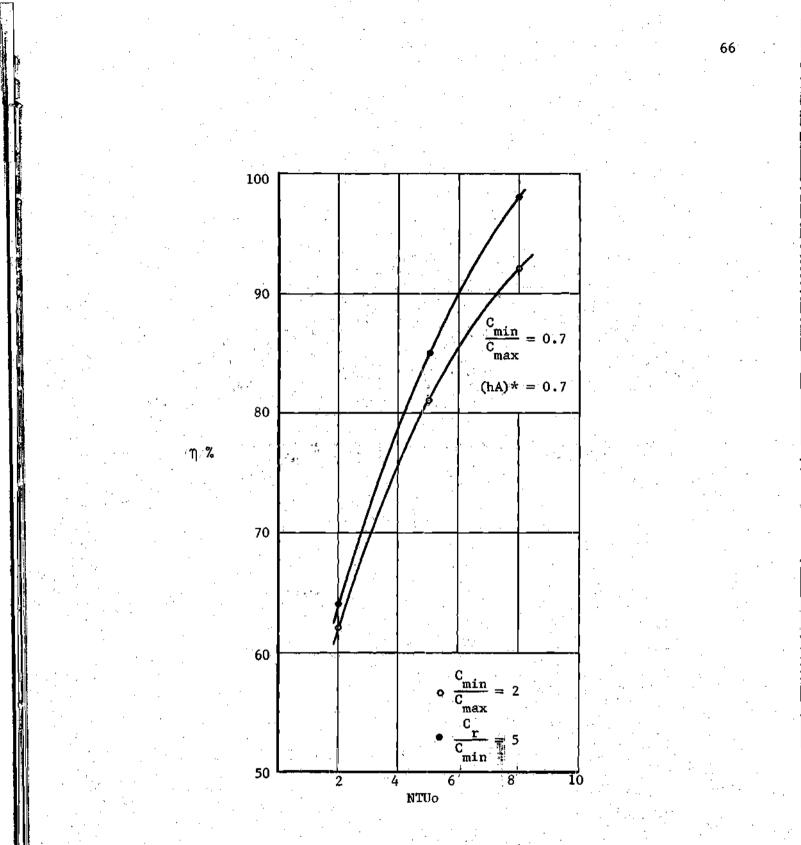
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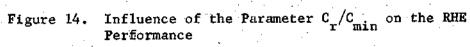
in the hot domain is reevaporated prior to reaching the cold side. Such is not the case for the lower average temperature distribution; rather, there is a carry-over of condensate to the cold side. This condensate being at a higher temperature than the temperature of the incoming cold air allows a greater heat transfer to the latter. Although this influence is limited to only the first column of elements on the cold side (due to immediate reevaporation of condensate due to heat transfer from the matrix), the end result is a slight increase in regenerator effectiveness.

Figure 14 shows the influence of C_r/C_{min} on the regenerator performance. It is seen that, at a higher value of this parameter, the effectiveness improves. This trend has been previously observed by other investigators and has been explained on the basis that, at a higher heat capacity rate of the matrix, its performance approaches that of a nonrotating conventional counterflow exchanger.

The influence of (hA)*, although reported by other investigators to be negligible for the range of variation considered in this study, seems to be slightly in variance with other studies, as shown in Figure 15. It is likely that this is more a consequence of computational errors involved with the present model and the resultant truncation error than it is an indication of a physical mechanism.

Figure 16 shows the influence of increasing the number of matrix elements on the effectiveness of the regenerator. That this influence is extremely small is a proper justification for extrapolating the results to the case of a very large number of elements.





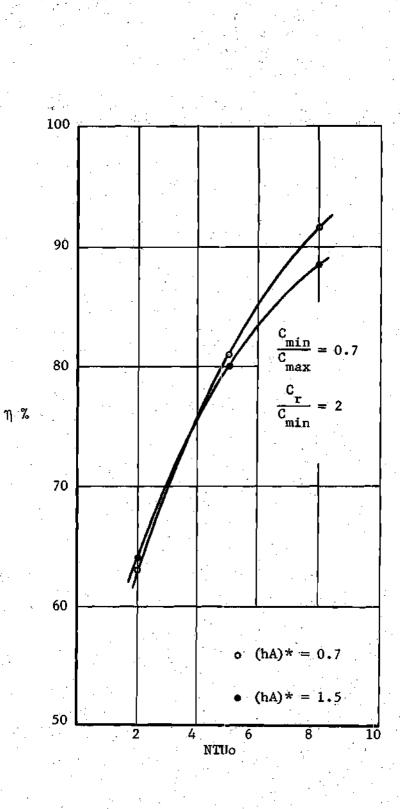
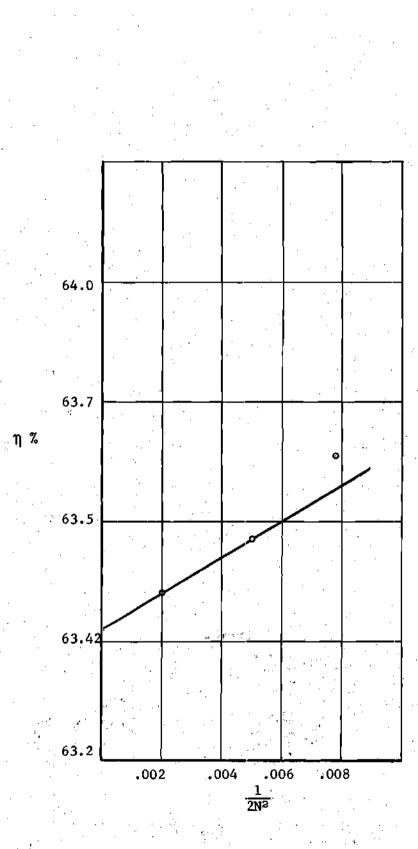
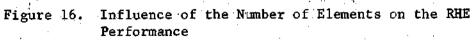


Figure 15. Influence of the Parameter (hA)* on the RHE Performance





CHAPTER VI

CONCLUSIONS AND RECOMMENDATIONS

The research reported in this investigation has demonstrated that i) It is feasible to examine rotary heat exchanger performance, with at least one condensing fluid stream through it, by the numerical technique developed in this work.

ii) Although an apparent increase in the rotary heat exchanger effectiveness occurs as a result of cold air humidification caused by a carry-over of condensate from the hot domain of the rotor matrix, the need to avoid humidification of cold air in certain applications puts an upper limit on the effectiveness of the heat exchanger.

For future investigations on the subject, it is suggested that i) Emphasis be placed on examining in more detail the influence of both the heat capacity rate parameters considered in this investigation, with particular reference to the quantitative definition of the condensate carry-over conditions:

if) For wider application of the numerical technique developed in this work, efforts be made to examine the influence of different entering gas stream states.

iii) A significantly large number of elements of the matrix be considered with truncation errors of computation limited to an order of magnitude smaller than those utilized in this work.

APPENDIX 1

DERIVATION OF EQUATIONS (3-27) AND (3-28)

Derivation of Equation (3-27) a)

. .

From Equation (3-23) and equation (3-20), we have

$$\dot{m}_{W}(i,j+1) = \frac{m_{a}}{N} \left[W(i,j) - W(i+i,j) \right] + \dot{m}_{W}(i,j).$$
 (A-1)

et
$$\dot{\mathbf{Q}} = \frac{\mathbf{m}}{\mathbf{N}} \left[\mathbf{h}_{\infty}(\mathbf{i},\mathbf{j}) - \mathbf{h}_{\infty}(\mathbf{i}+1,\mathbf{j}) \right]$$
 (A-2)

Incorporating Equations (A-1) and (A-2) into Equation (3-21) yields $\sum_{i=1}^{n-1} \frac{\sigma_i \sigma_i}{\sigma_i \sigma_i} > 0.4$

$$\dot{Q} = \frac{m_{a}}{N} \left[W(i,j) + W(i+1,j) \right] C_{pw} T(i,j+1) + \left[m_{w}(i,j) C_{pw} + (A-3) \right]$$

$$\dot{M}_{r} C_{pr} \left[T(i,j+1) - T(i,j) \right] - \frac{kA_{d}}{2L} \left[(T(i-1,j) + T(i-1,j+1)) - (T(i,j) - T(i,j)) \right]$$

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$$2(T(i,j) + T(i,j+1)) + (T(i+1,j) + T(i+1,j+1))$$

From Equation (3-24)

$$W(i+1,j) = \begin{bmatrix} \frac{i}{m_{a}} - \frac{fAv}{2C} \\ \frac{pm}{m_{a}} + \frac{fAv}{2C} \end{bmatrix} W(i,j) + \begin{bmatrix} \frac{fA}{2C} \\ \frac{pm}{m_{a}} + \frac{fAv}{2C} \end{bmatrix} \begin{bmatrix} (A-4) \\ 2W_{savg}(i,j) \end{bmatrix}$$

For simplification let

$$\frac{fA}{2m} \frac{C}{C} \frac{N}{pm} = \frac{B}{N} = \frac{NTU}{2N}$$
 (A-5)

Equation (A-4) can now be rewritten as

$$W(i+1,j) = \frac{N-B}{N+B} W(i,j) + \frac{2B}{N+B} W_{savg}(i,j)$$
 (A-6)

Incorporation of Equation (A-6) into Equation (A-3) yields

~

$$= \frac{m_a}{N} \left[\left(1 - \frac{N-B}{N+B} \right) W(i,j) - \frac{2B}{N+B} W_{savg}(i,j) \right] C_{pw} T(i,j+1) +$$

$$\left[m_w(i,j) C_{pw} + \frac{m_r C_{pr}}{N} \right] T(i,j+1) - T(i,j) - T(i,j)$$

 m_a HC84 $\lfloor (T(i-1,j) + T(i-1,j+1)) - 2(T(i,j) + T(i,j+1)) +$

(T(i+1,j) + T(i+1,j+1)),

where $HC84 = \frac{kA}{2m_{a}L}$

Finally combining Equations (A-7) and (A-8), one obtains

$$h_{\infty}(i+1,j) = h_{\infty}(i,j) - \frac{2B}{N+B} \left[W(i,j) - W_{savg} \right] C_{pw} T(i,j+1) - (A-9)$$

(Continued)

(A-8)

$$\left[\frac{\stackrel{m}{w}(i,j)}{\stackrel{c}{\frac{m}{pw}}}_{\frac{m}{N}} + \frac{\stackrel{m}{r} \stackrel{c}{\frac{pr}{p}}_{a}}{\stackrel{m}{a}}\right] T(i,j+1) = T(i,j) - HC84 N \left[(T(i-1,j) + \frac{m}{N}) \right]$$

T(i-1, j+1)) - 2(T(i, j) + T(i, j+1)) + (T(i+1, j) + T(i+1, j+1)).

If we incorporate Equations (3-29), (3-30), and (3-31) into Equation (A-9), we will have the resulting Equation (3-27).

b) Derivation of Equation (3-28)

From Equation (3-22) we have

$$h_{\infty}(i+1,j) = \frac{N-B}{N+B} h_{\infty}(i,j) + \frac{2B}{N+B} h_{savg}(i,j)$$

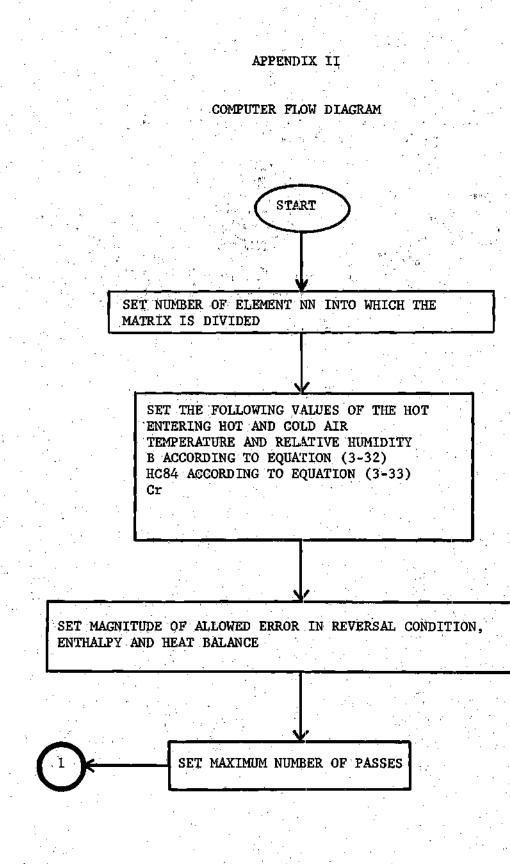
Adding and subtracting $\frac{2B}{N+B} h_{\infty}(1,j)$,

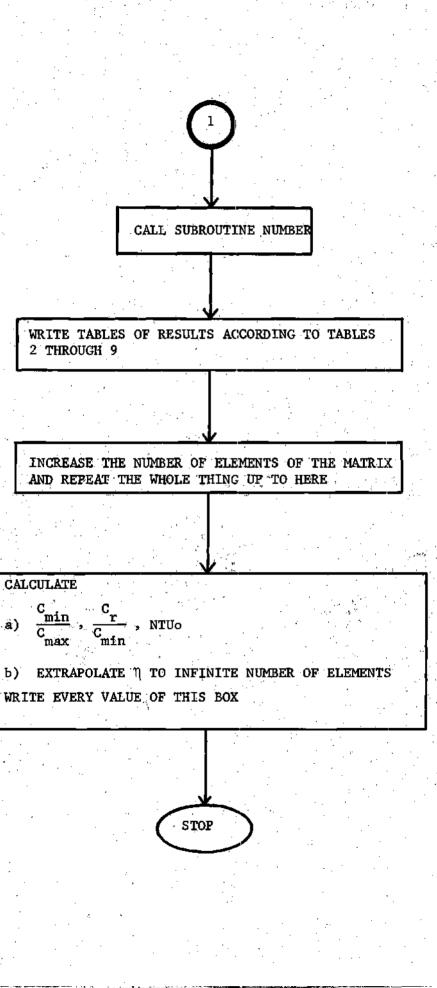
 $h_{\omega}(i+1,j) = h_{\omega}(i,j) + HC81 (h_{savg}(i,j) - h_{\omega}(i,j)),$

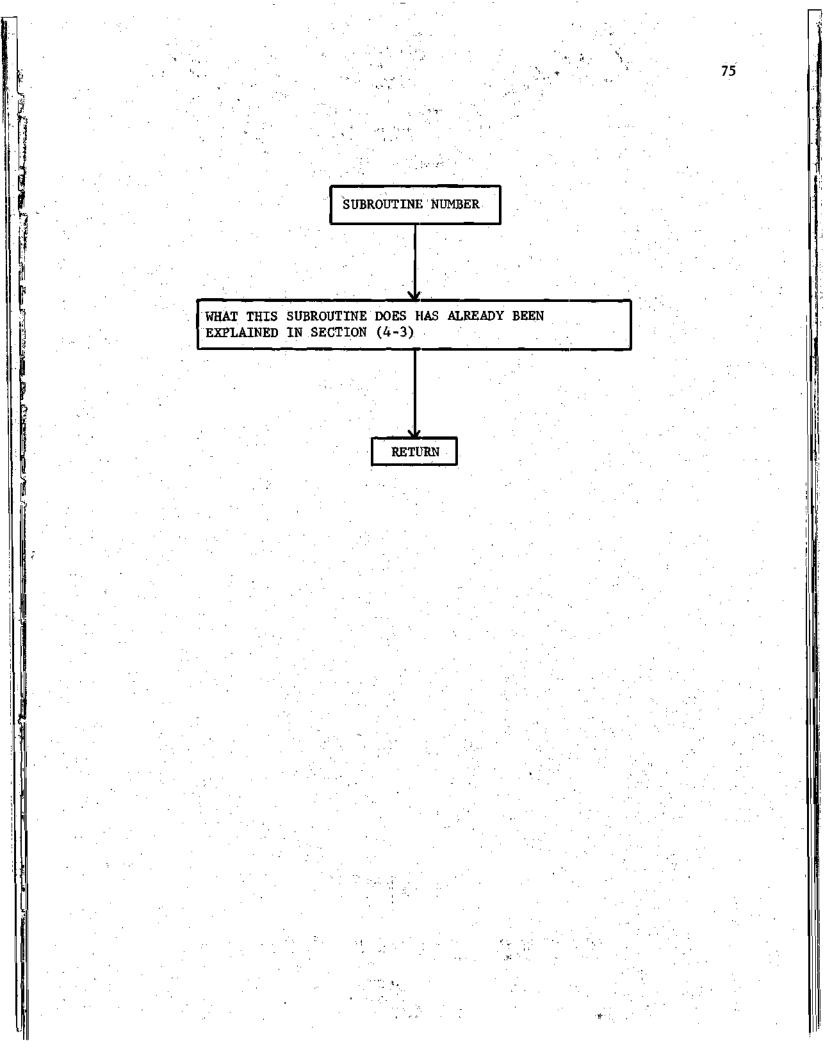
which is the same as Equation (3-28).

(A-10)

(A-11)







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